

Active Galactic Nuclei: X-ray surveys and AGN evolution

On the attempt to “replace” the Unified Model for AGN
by the AGN/galaxy co-evolution prescriptions

Two main themes in modern high-energy astrophysics

□ Physics of accretion and ejection in massive black holes

Needs characterization of the X-ray and γ -ray emission from AGN, hence high counting statistics (large effective area) and, possibly, high-resolution X-ray spectra. [Lessons by Dr.ssa P. Grandi/E. Torresi and Dr. M. Cappi]

Census of SMBHs to “map” the growth of massive structures up to high redshifts: AGN/galaxy co-evolution, feedback processes, etc.

Needs large, well-defined samples of AGN, including the most elusive, heavily obscured ones, and the first SMBHs to form in the Universe. Large source numbers are more important than individual source photon statistics, typically very limited (e.g., in deep X-ray surveys).

Outline

- ✓ AGN Unified scheme vs. AGN/galaxy co-evolution models
- ✓ The first massive black holes
- ✓ Integrated AGN emission recorded in the X-ray background (XRB) and the role of obscured AGN
- ✓ X-ray surveys: depth vs. coverage
- ✓ New insights into the X-ray absorber (torus) from mid-IR observations

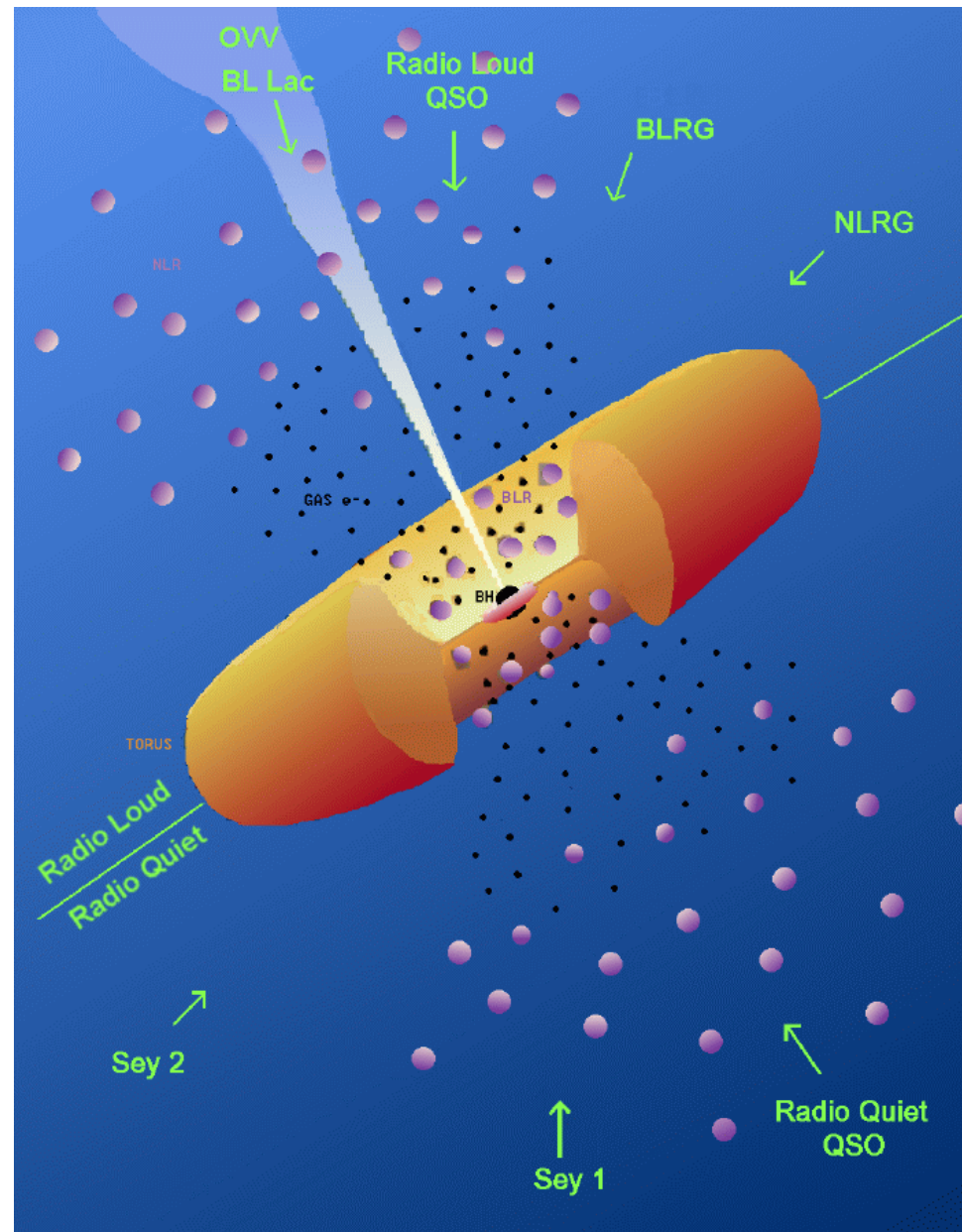
For a recent review on the subject, see Alexander & Hickox 2012, *New Astronomy Reviews*, 56, p. 93 (arXiv:1112:1949) – see also Brandt & Alexander 2015, *The Astronomy & Astrophysics Review*, 23, p.93 (arXiv:1501.01982)

AGN Unified Model

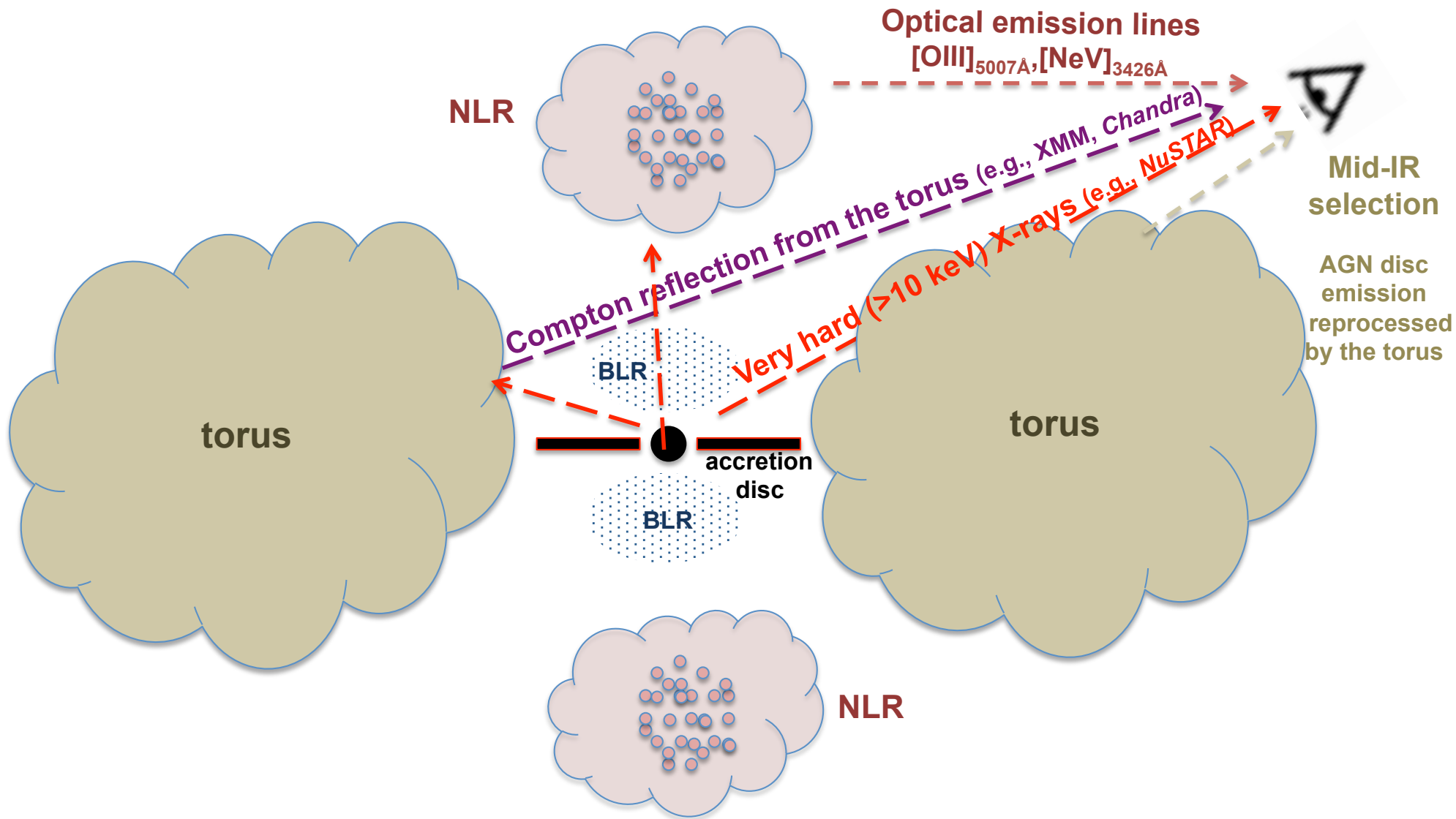
after Antonucci & Miller 1985;
Antonucci 1993

Fine for many AGN as
a baseline for the
description of different
observational
properties

Probably not the end
of the story



adapted from Urry
& Padovani 1995



A logarithmic view of an AGN

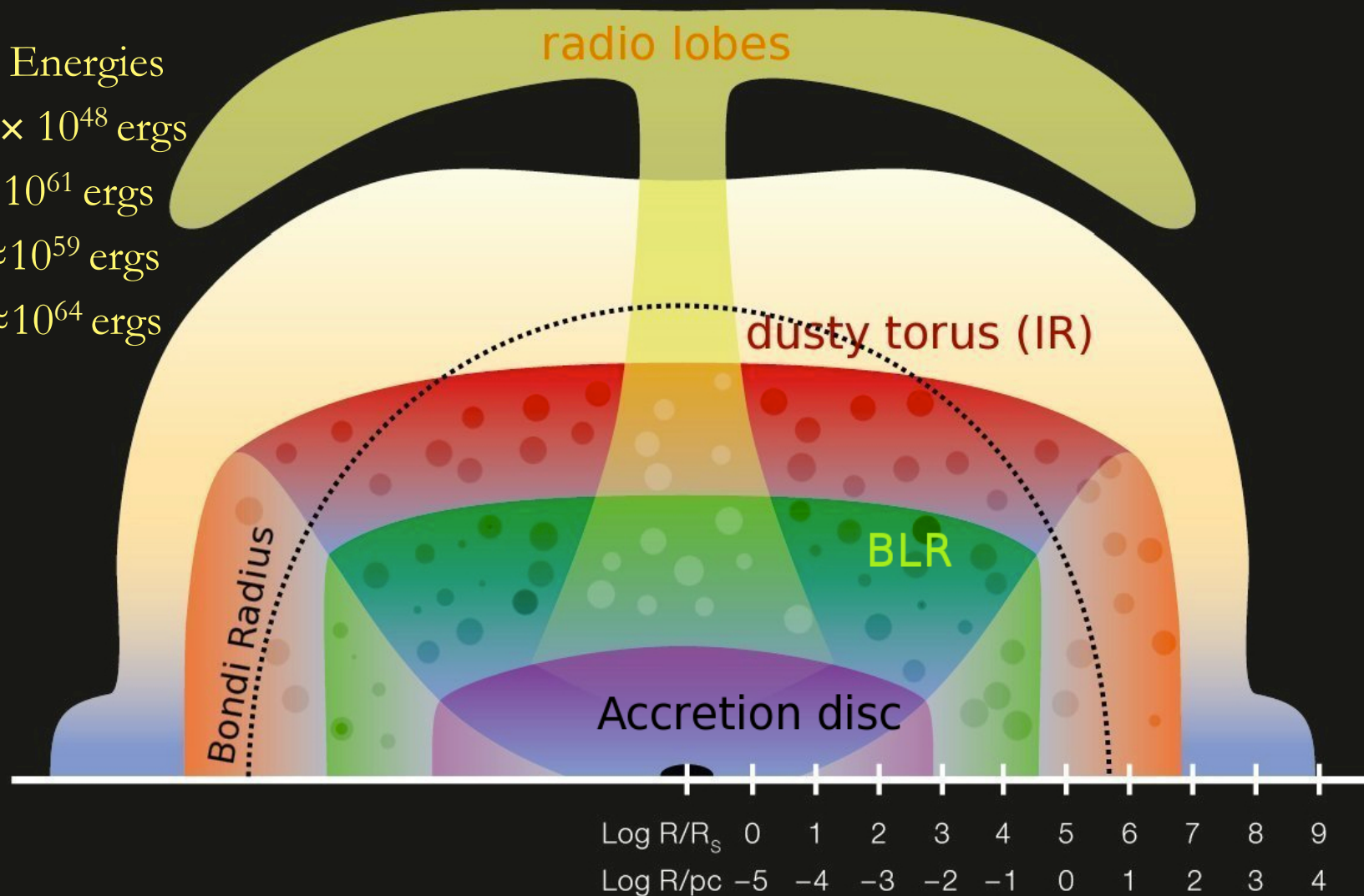
Binding Energies

$$E_{b,\odot} \approx 4 \times 10^{48} \text{ ergs}$$

$$E_{b,\text{BH},8} \approx 10^{61} \text{ ergs}$$

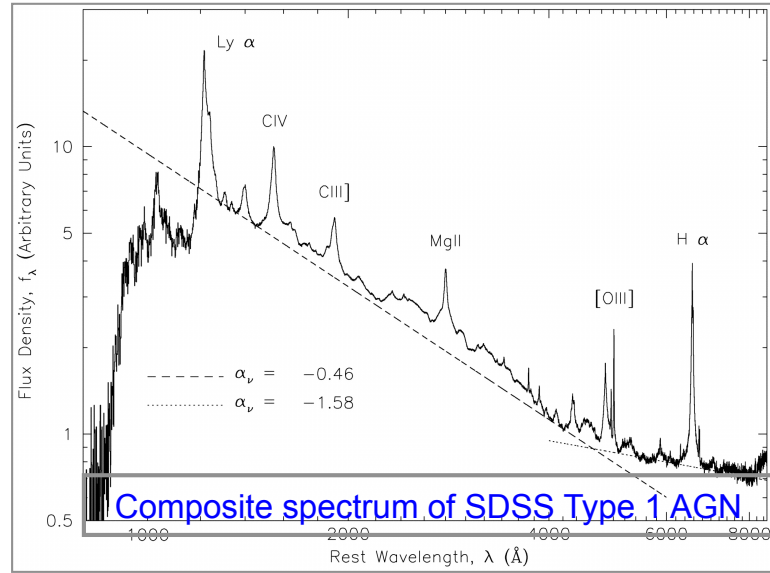
$$E_{b,\text{gal},11} \approx 10^{59} \text{ ergs}$$

$$E_{b,\text{Coma}} \approx 10^{64} \text{ ergs}$$

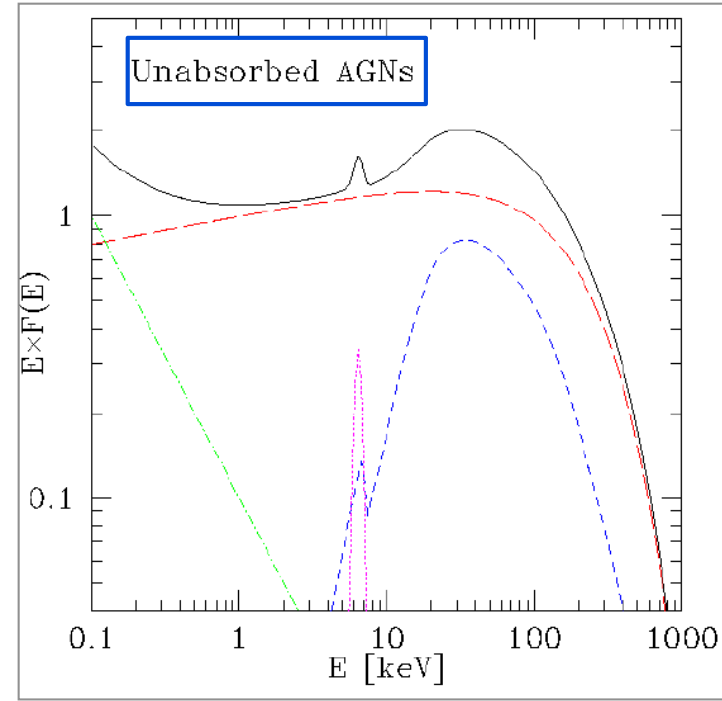


Courtesy of A. Merloni, ESO graphics, 2010

Optical band

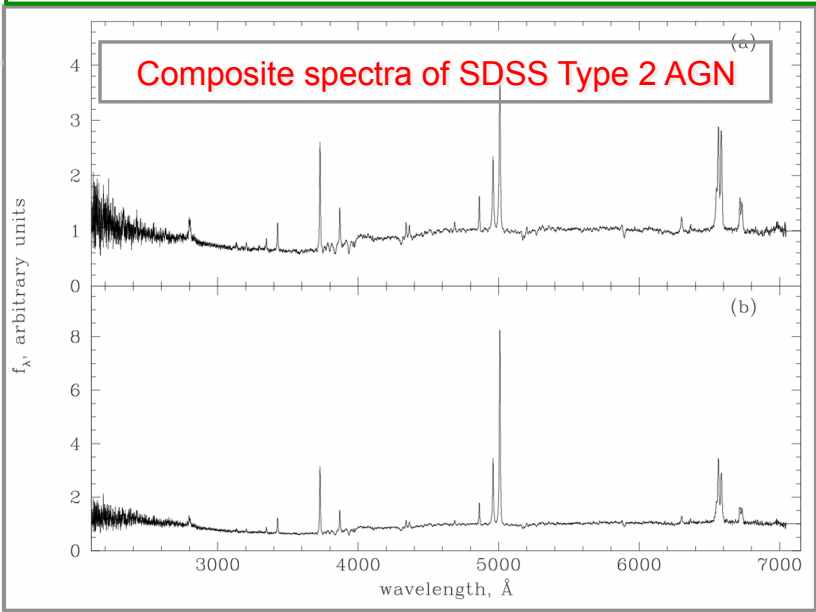


Type 1
AGN

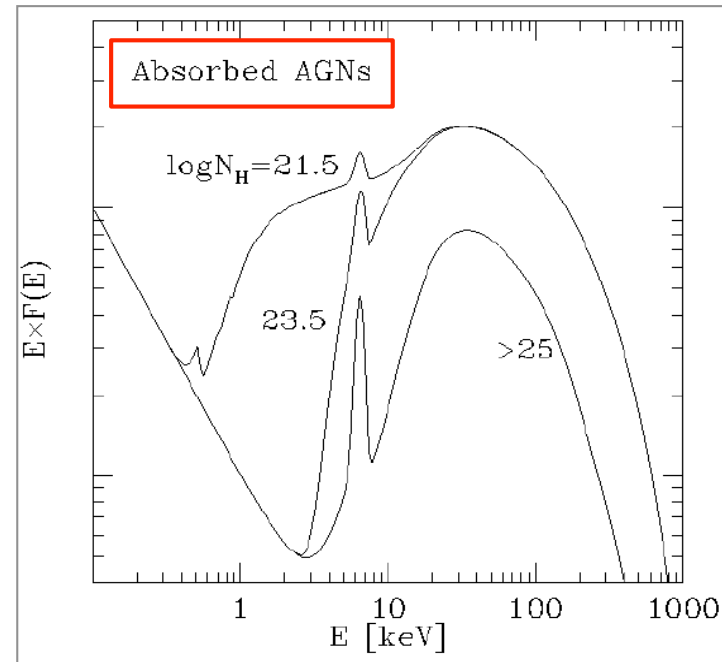


X-ray band

Type 2 AGN easily missed in optical and partly in X-ray surveys

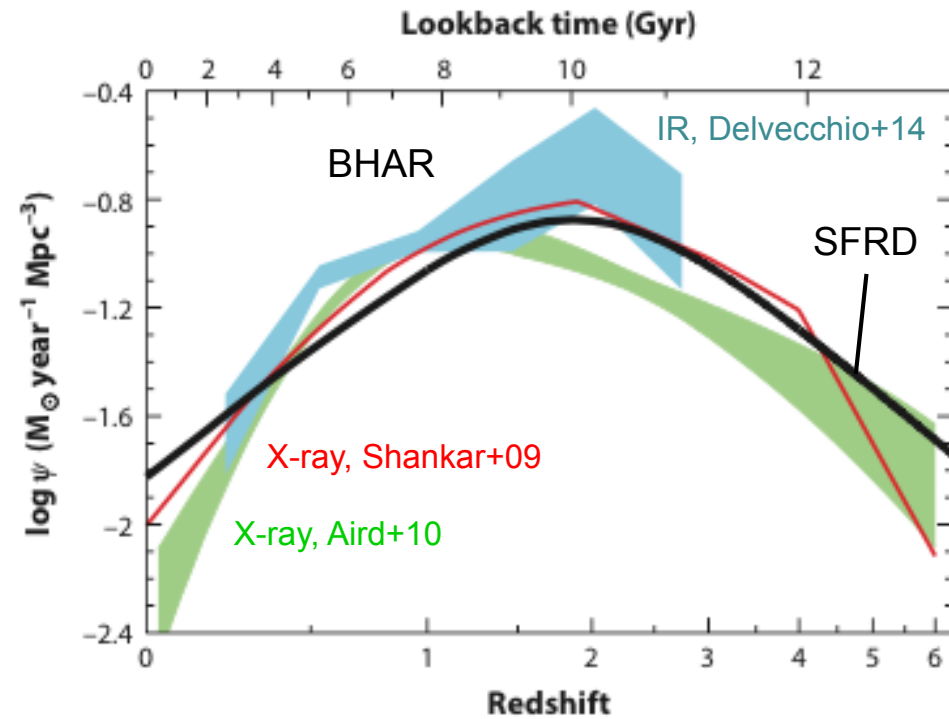
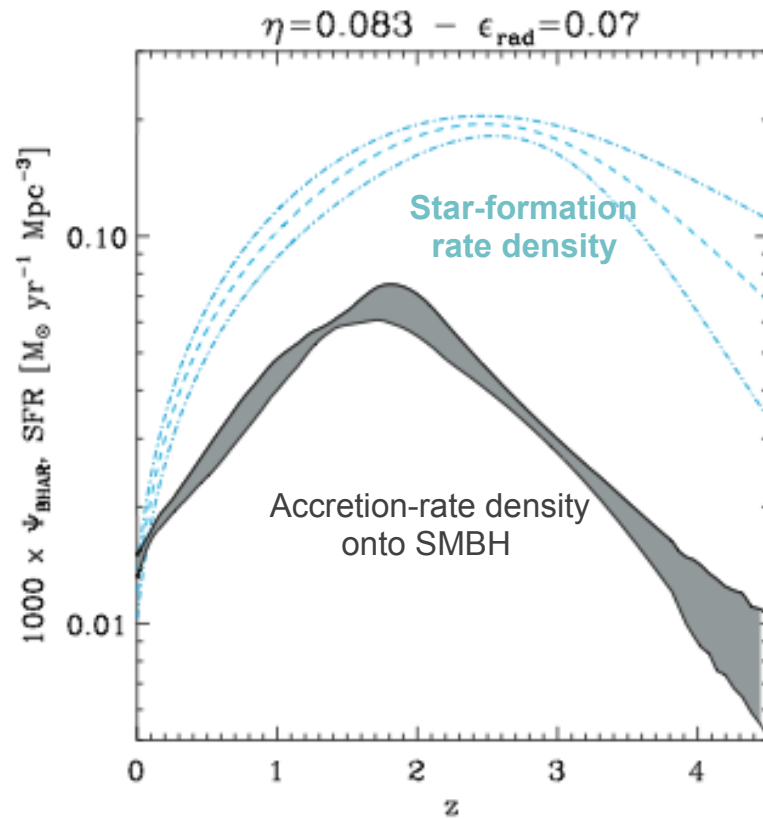


Type 2
AGN



AGN-galaxy co-evolution

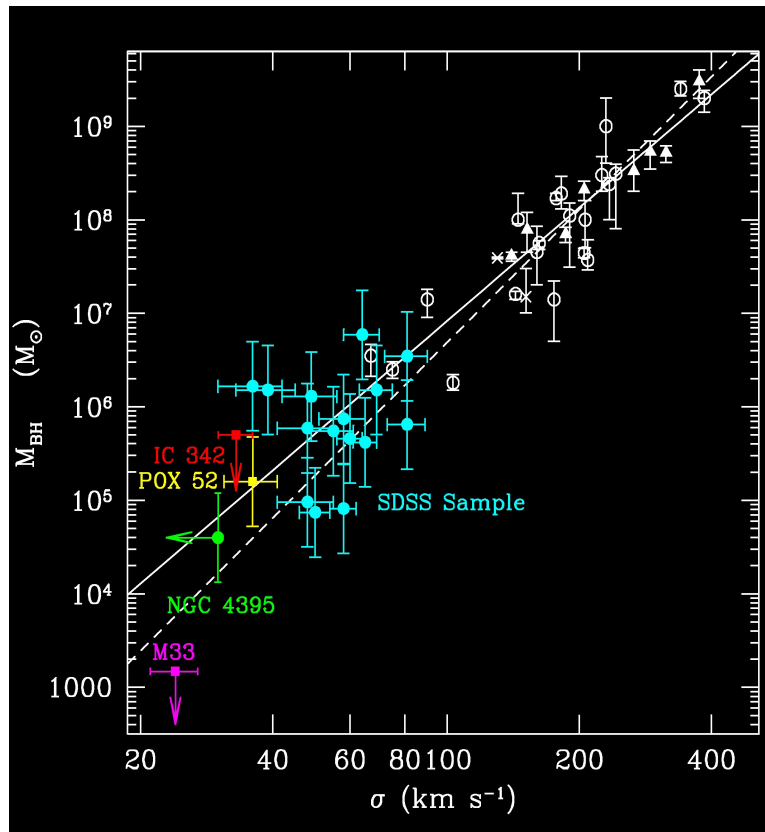
Accretion and star formation over cosmic



from Merloni & Heinz 2008;
see also Hopkins & Beacom 2006, Gruppioni et al. 2011

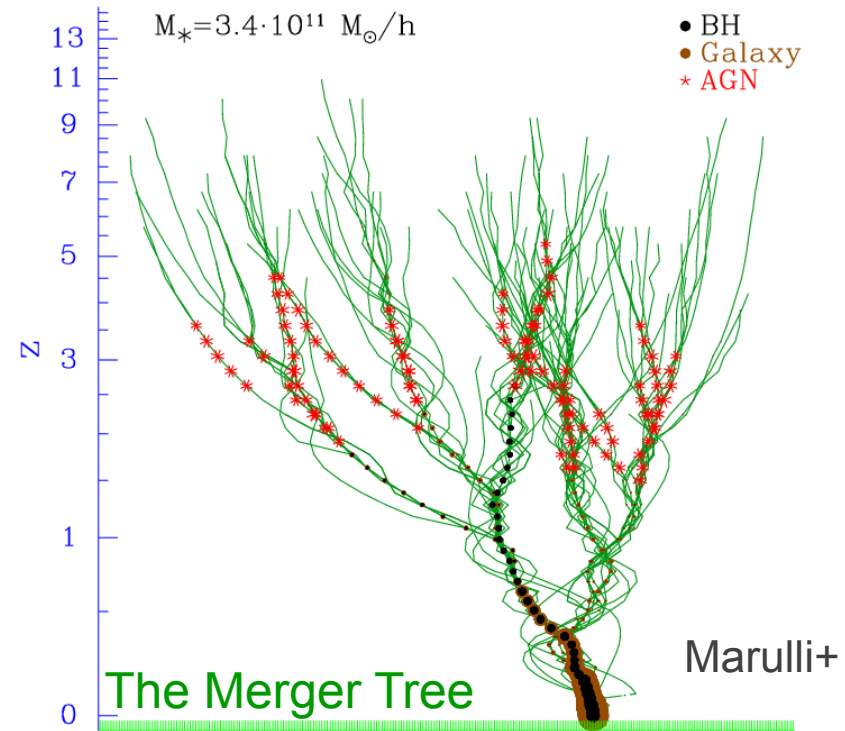
from Madau & Dickinson 2014

AGN as a key phase of a galaxy lifetime



Scaling relations between BH mass
and host galaxy properties
(stellar bulge mass, luminosity,
velocity dispersion)

AGN and galaxies closely tied
→ co-evolution

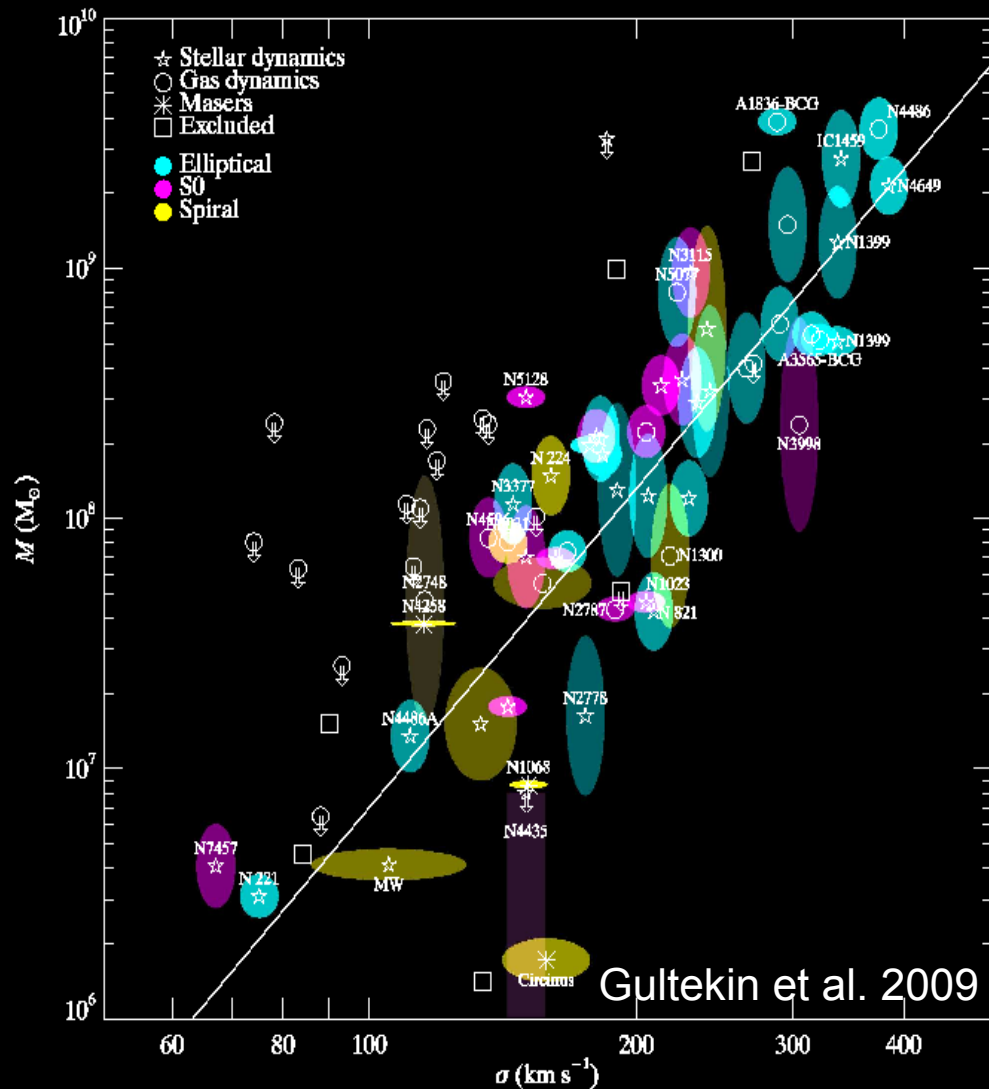


**Semi-analytic models of BH/galaxy
co-evolution** (e.g: Kauffmann+98, Volonteri
+06, Salvaterra+06, Rhook&Haehnelt08,
Hopkins+08, Menci+08, Marulli+09)

These follow the evolution and merging of Dark
Matter Halos with cosmic time and use
analytic recipes to treat baryon physics.

Condition: nuclear trigger at merging

Black Hole – galaxy scaling relations



Correlation between BH mass and galaxy velocity dispersion σ

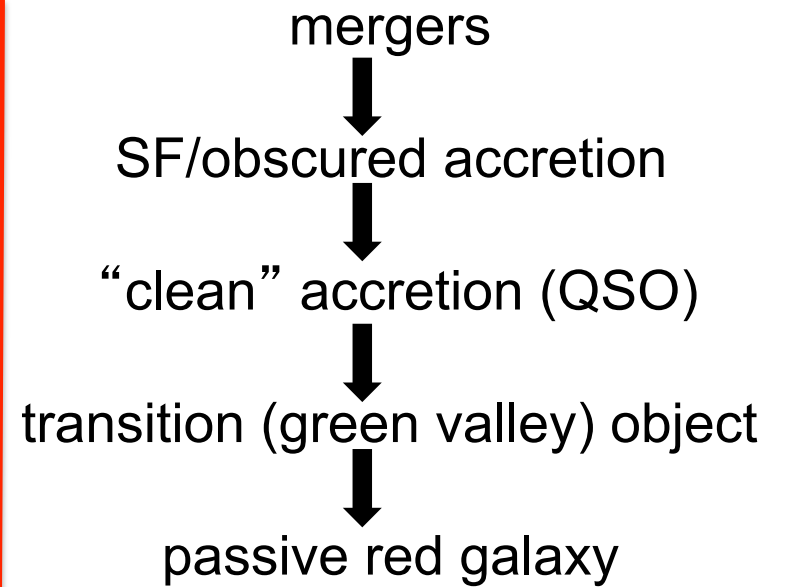
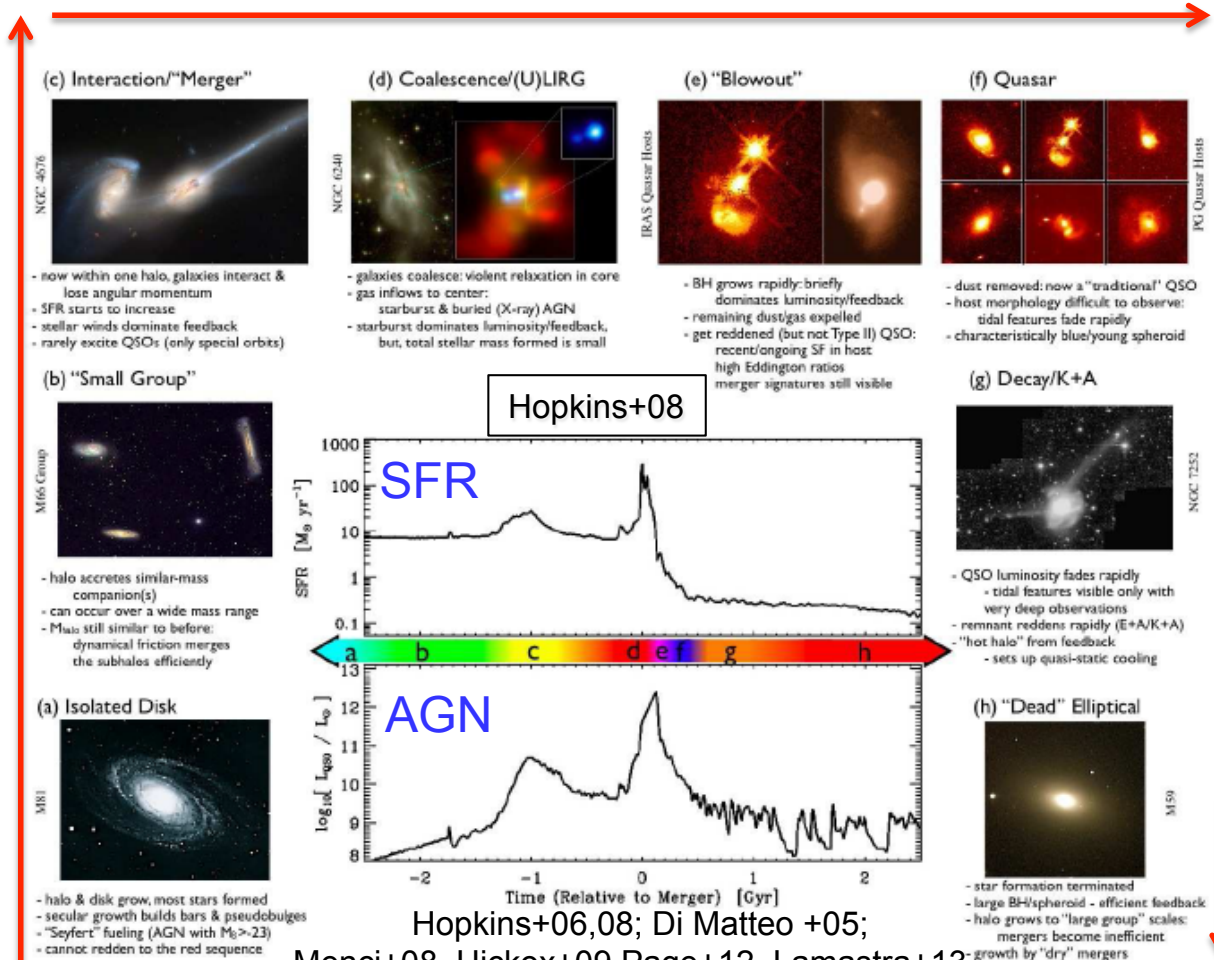
σ measured well **outside** the gravitational sphere of influence of the BH

No causal connection (now)

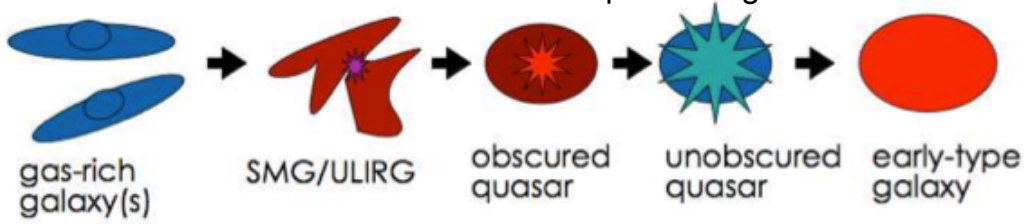
Either coincidence (!) or the result of **common evolution**

Kormendy and Richstone 1995; Magorrian et al. 1998; Gebhardt et al. 2000; Ferrarese et al. 2000; Tremaine et al. 2002; Gultekin et al. 2009; Kormendy & Bender 2012 – see also Jahnke & Maccio' 2011

The BH/galaxy “evolutionary sequence”



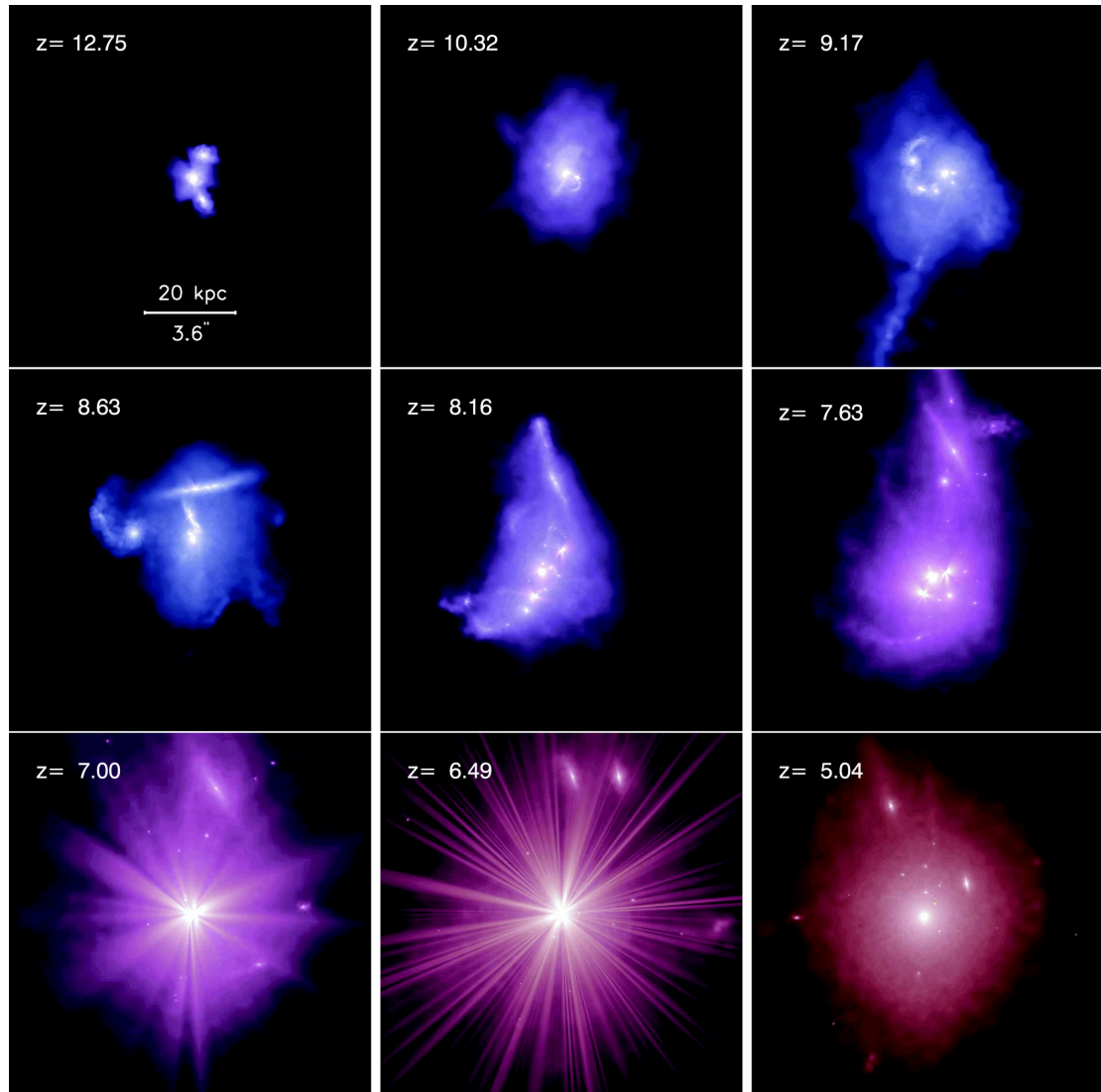
Importance of merger-driven vs. secular (“smooth”) accretion (e.g., Elbaz+11, Rodighiero+11, Rovilos+12)



Strong winds (=feedback) expected in the “blowout” phase

Hickox+09

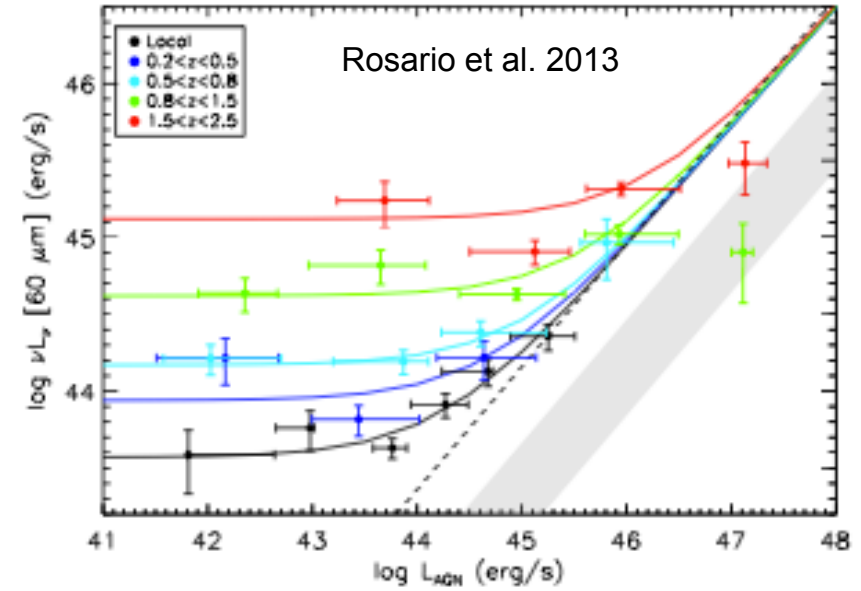
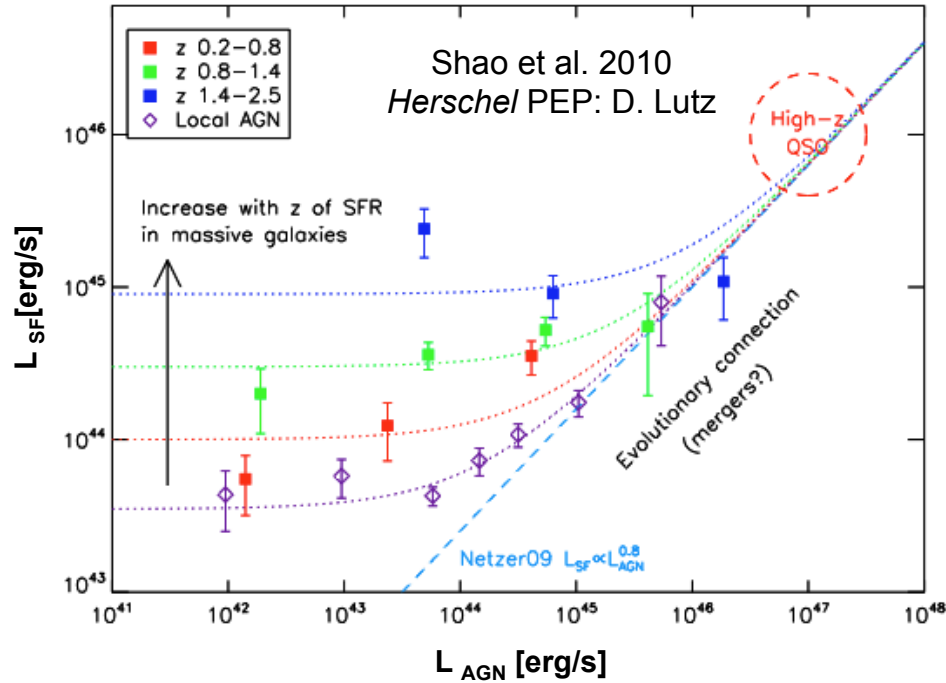
Simulated formation of a $\approx 10^9 M_{\odot}$ BH at high z



- Early on
 - strong galaxy interactions= violent star-bursts
 - heavily obscured QSOs
- When galaxies coalesce
 - accretion peaks
 - QSO becomes optically visible as AGN winds blow out gas
 - outflows as direct evidence for strict QSO/galaxy relation (feedback)
- Later times
 - SF & accretion quenched
 - red spheroid, passive evolution

Li+07

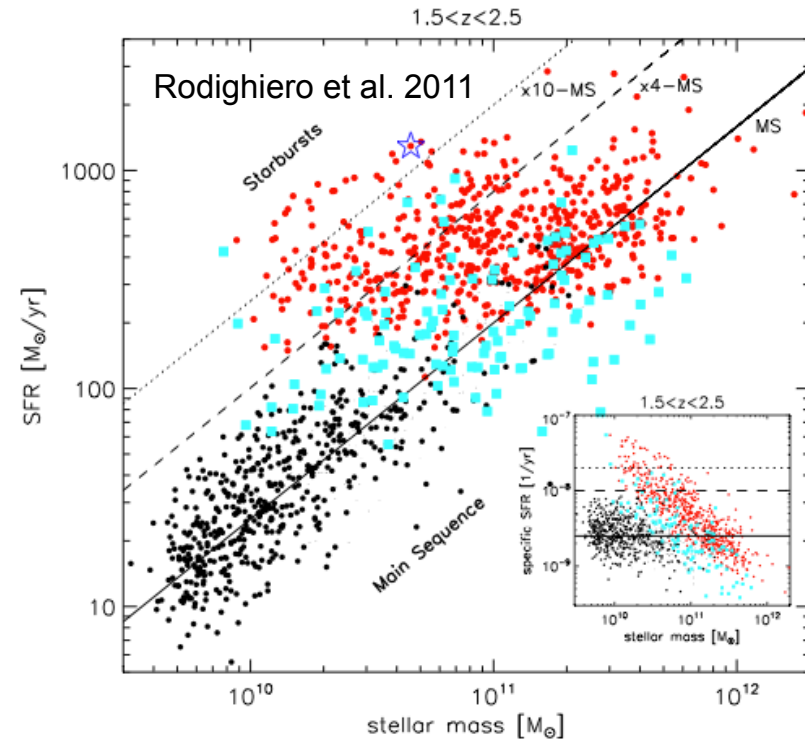
[$M_{\text{BH}} - \sigma - M_{\text{Bulge}} - \dots$ relations]



Two paths of AGN/galaxy co-evolution

- At high AGN luminosity, galaxy merging is the driver of accretion and star formation → rapid bursts of activity (~10% population?)
- At lower AGN luminosity, SF has little dependence on AGN luminosity → secular, non-merger driven star formation (~90% pop?)

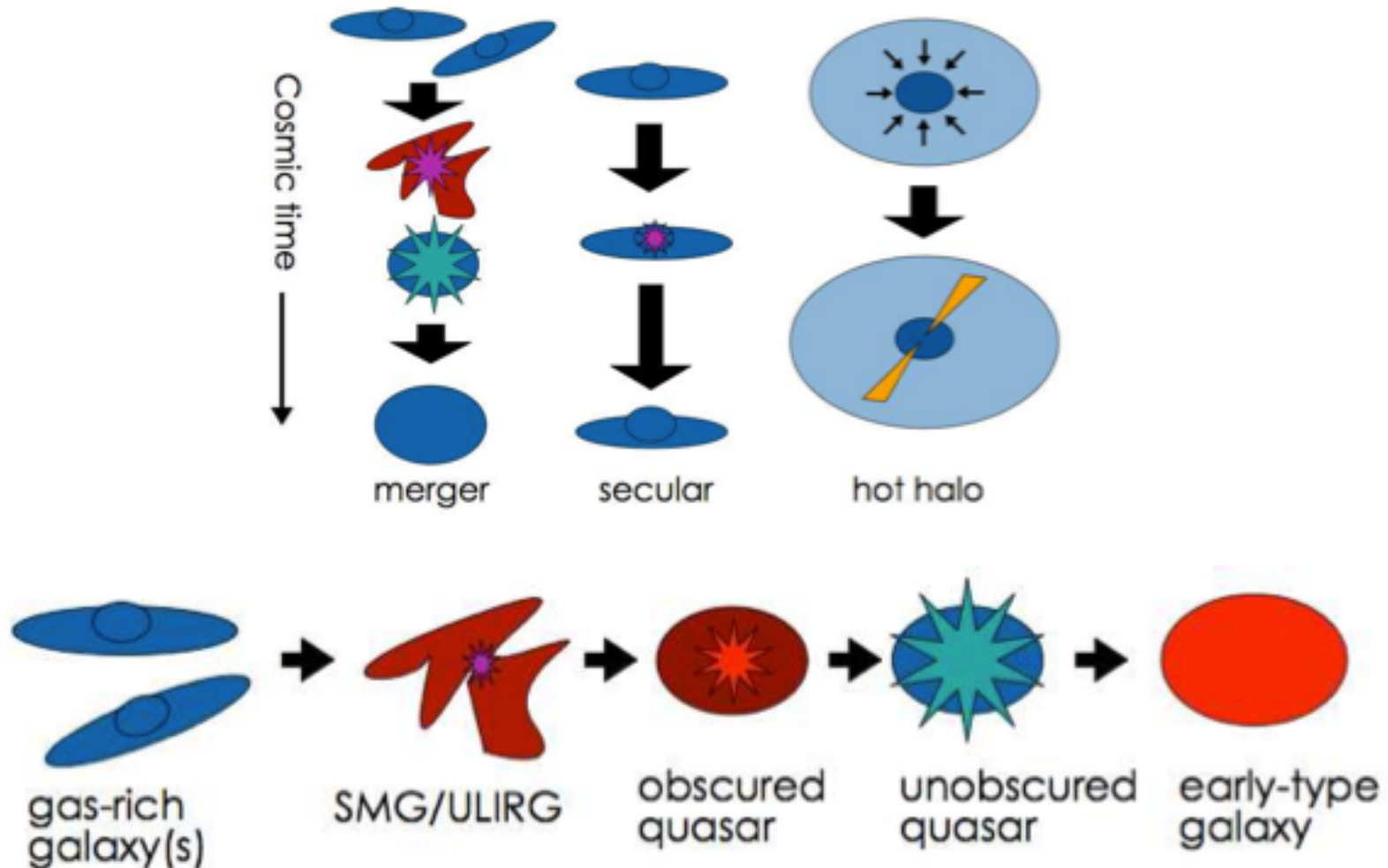
(e.g. Georgakakis+09, Lutz+10, Cisternas+11, Schawinski+11, Elbaz+11, Rodighiero+11, Mullaney+11, Santini+11, Rovilos+12, Rosario+12, ...)



Two modes of accretion:

Mergers \leftrightarrow luminous quasars

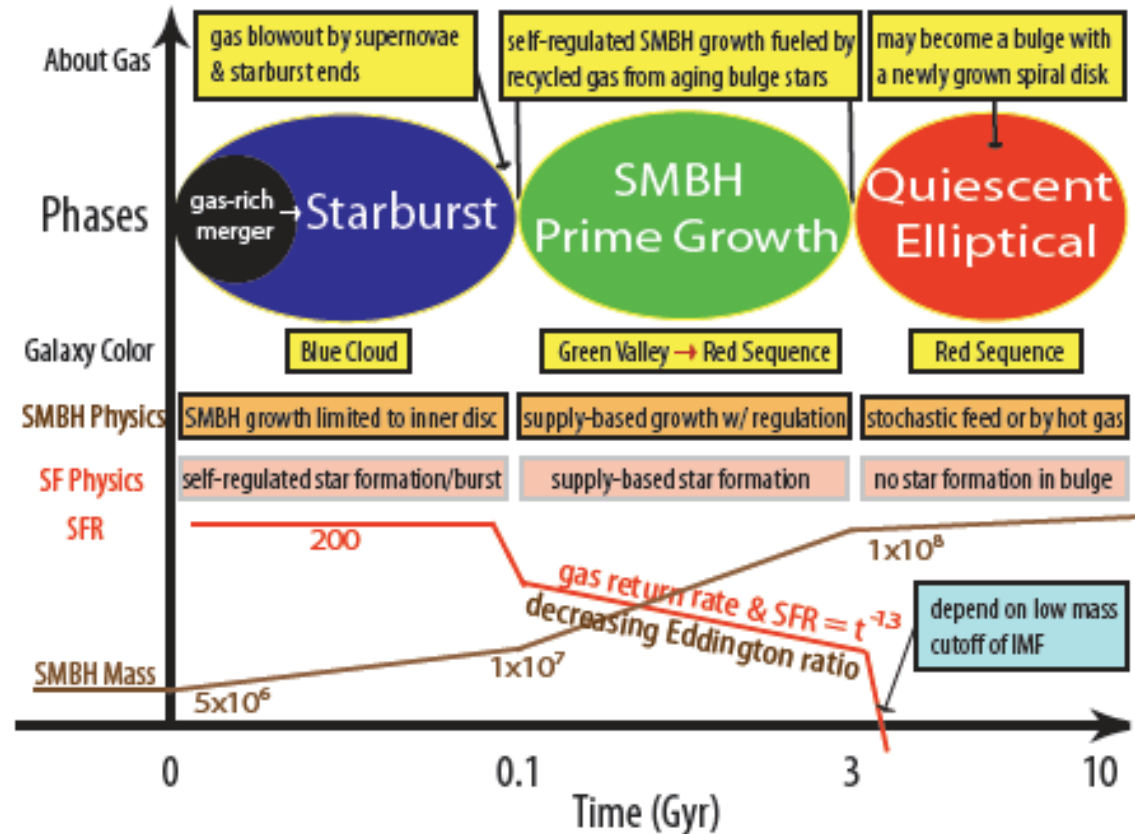
Secular (disk instabilities, bars, minor mergers) \leftrightarrow low-luminosity AGN

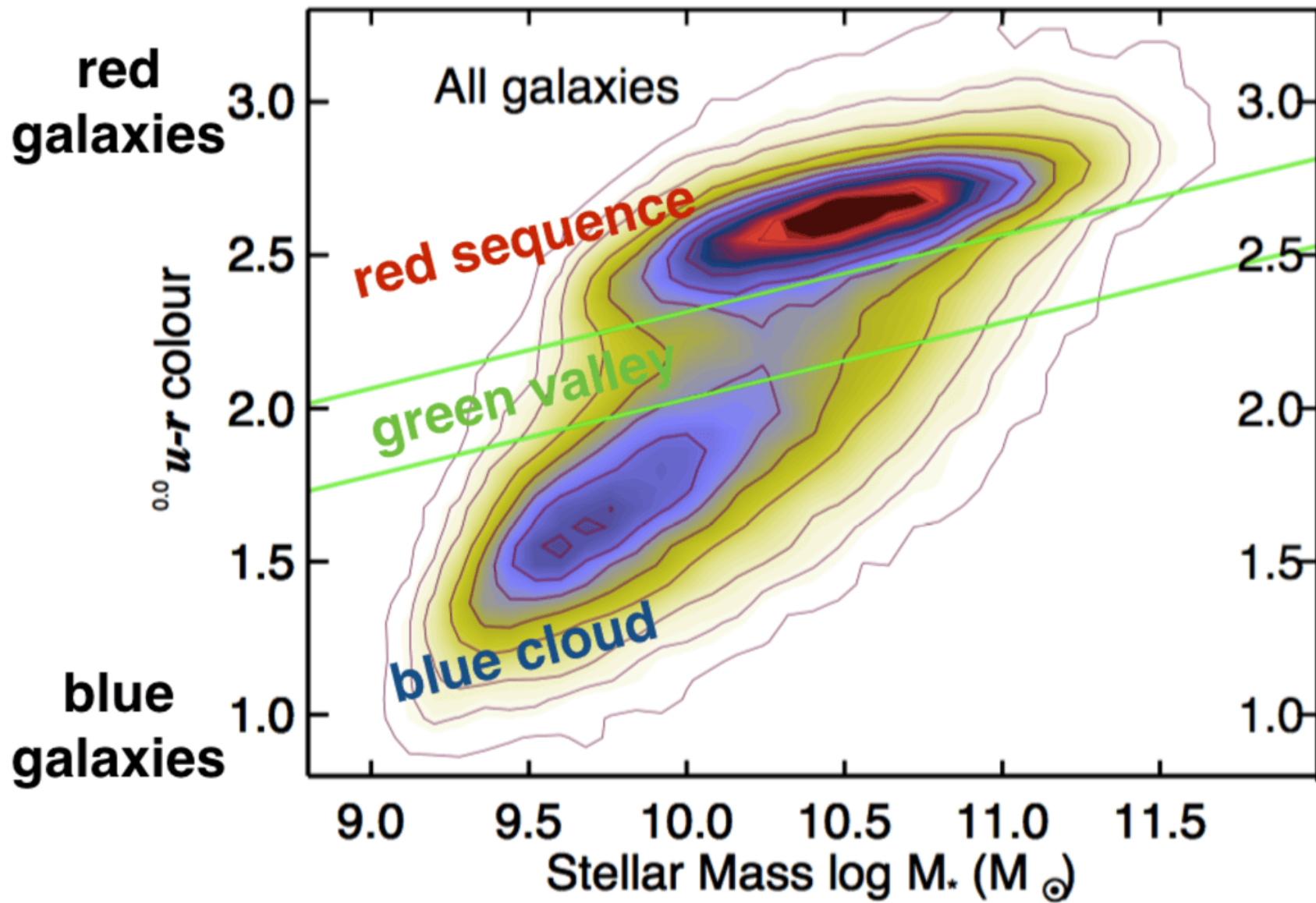


An alternative picture

STB precedes SMBH growth, lasts 10-100 Myr, and then stops itself (through SN)

Main SMBH growth in the post-starburst phase fueled by recycled gas from inner bulge (old) stars and lasts $\gg 100$ Myr, albeit at relatively low and diminishing Eddington ratios for most of the time





Claims of a higher fraction of AGN activity in the green valley: higher availability of fuel?

Obscured AGN growth and star formation at $z \approx 2$

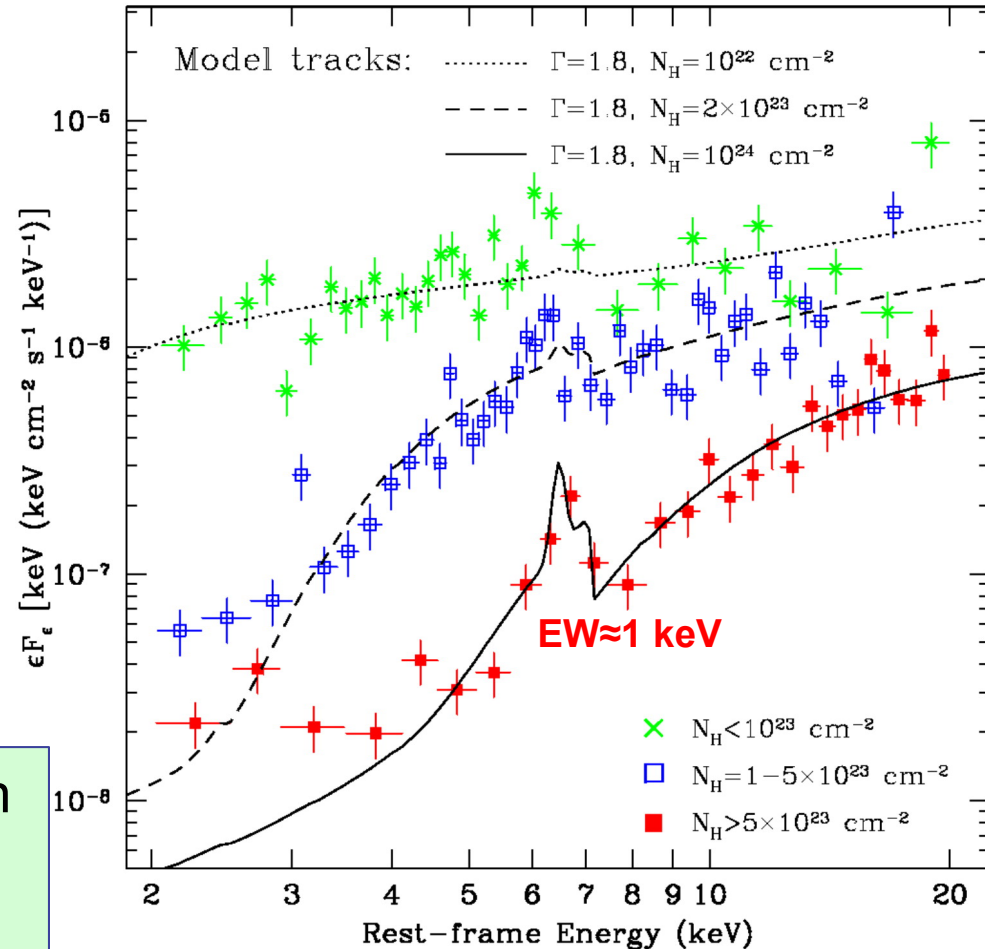
Obscured AGN in sub-mm galaxies

Large reservoir of gas available for accretion and SF

Further indications from mid-IR/optical selected sources

Deep X-ray fields and stacking techniques needed to estimate average source properties

Obscured accretion = key phase in AGN growth and AGN/galaxy co-evolution → Much of the mass growth of SMBH occurs during the heavily obscured phase (e.g., Treister+10)



Alexander et al. 2005

→ Needed: census and knowledge of Compton-thick AGN

Two main open issues

High-redshift

BH/galaxy co-evolution
still unconstrained at very
high- z ($z > 6$ or so).
Already formed luminous
QSOs at $z=6$

Heavily obscured AGN

Heavily obscured accretion
mostly unconstrained
beyond the local Universe



Requirement: a complete census of AGN activity

Information stored in the X-ray background

Open issue: time for BH growth at $z \approx 6$

Growth of BHs: trade-off between the gas “converted” into radiation and that accreted onto the SMBH

$$M(t) = M_0 \exp\left(\frac{1 - \epsilon}{\epsilon} \frac{t}{t_{\text{Edd}}}\right)$$

Larger radiation efficiency ϵ means longer times to achieve a given mass

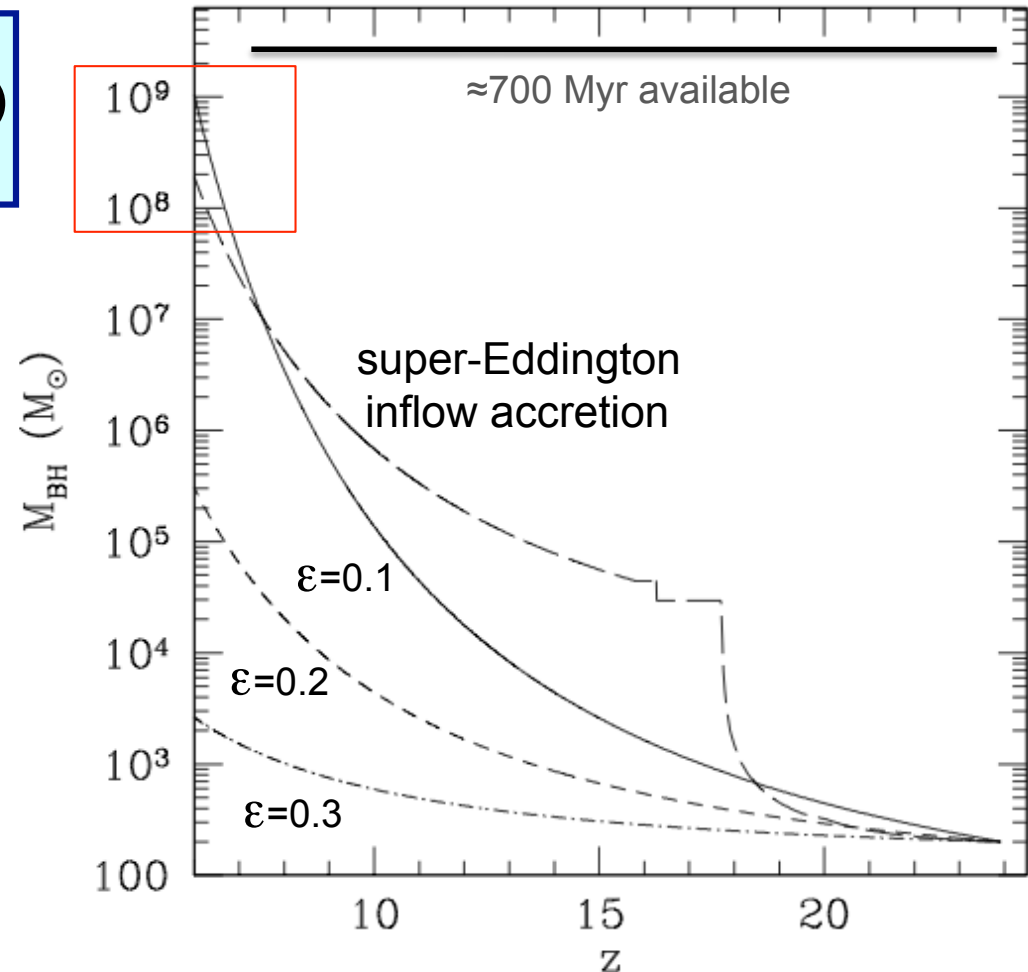
[$t_{\text{Edd}} = 0.45$ Gyr for $\epsilon = 0.1$]

Rapidly spinning BHs might have problems because of a larger ϵ

Highest-redshift quasar so far spectroscopically identified:

ULASJ1120+0641, $z = 7.08$,

$M_{\text{BH}} \approx 2 \times 10^9 M_{\odot}$ (Mortlock et al. 2011)

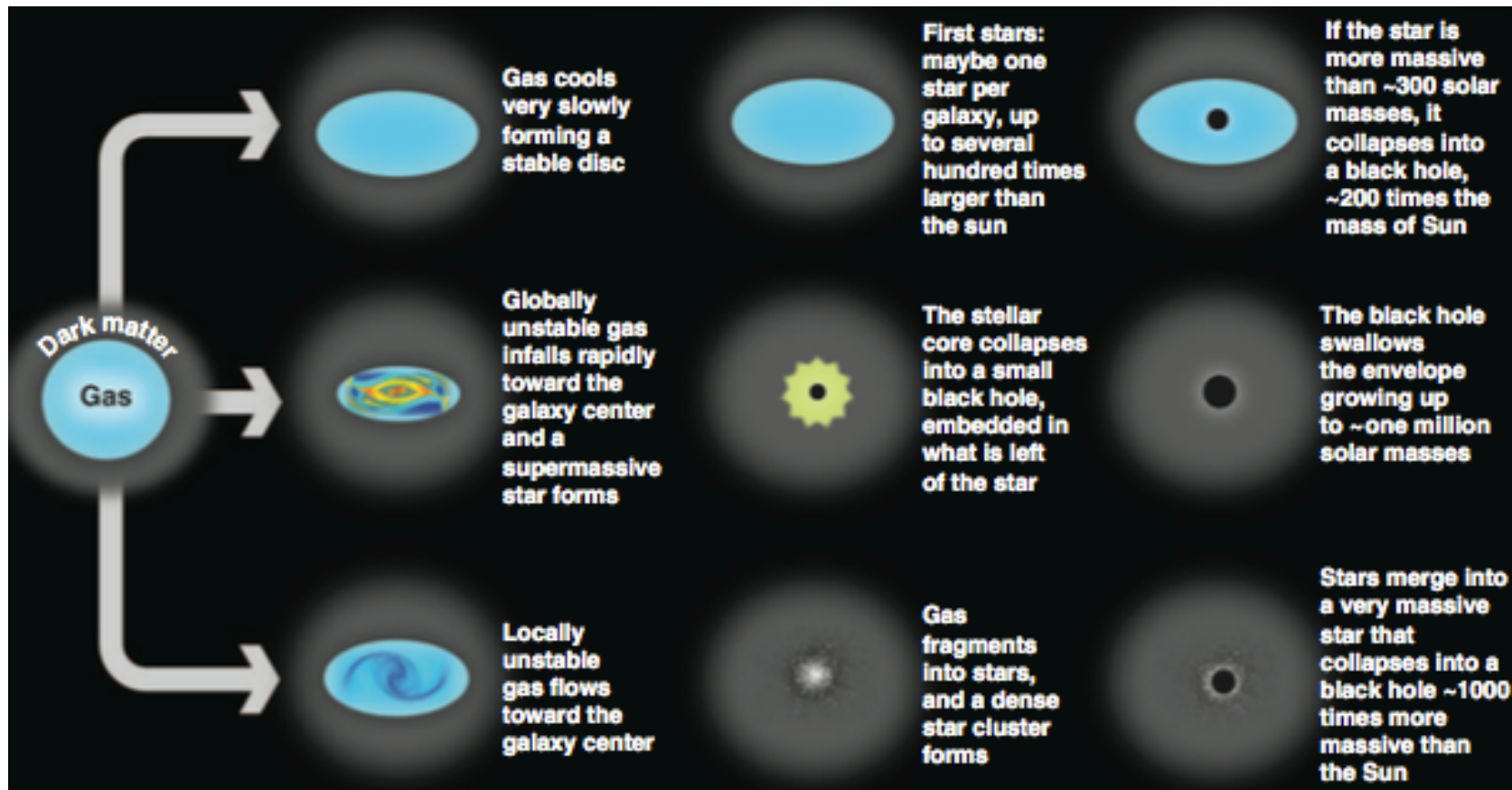


Possible problems with the mass of the “seed” BHs

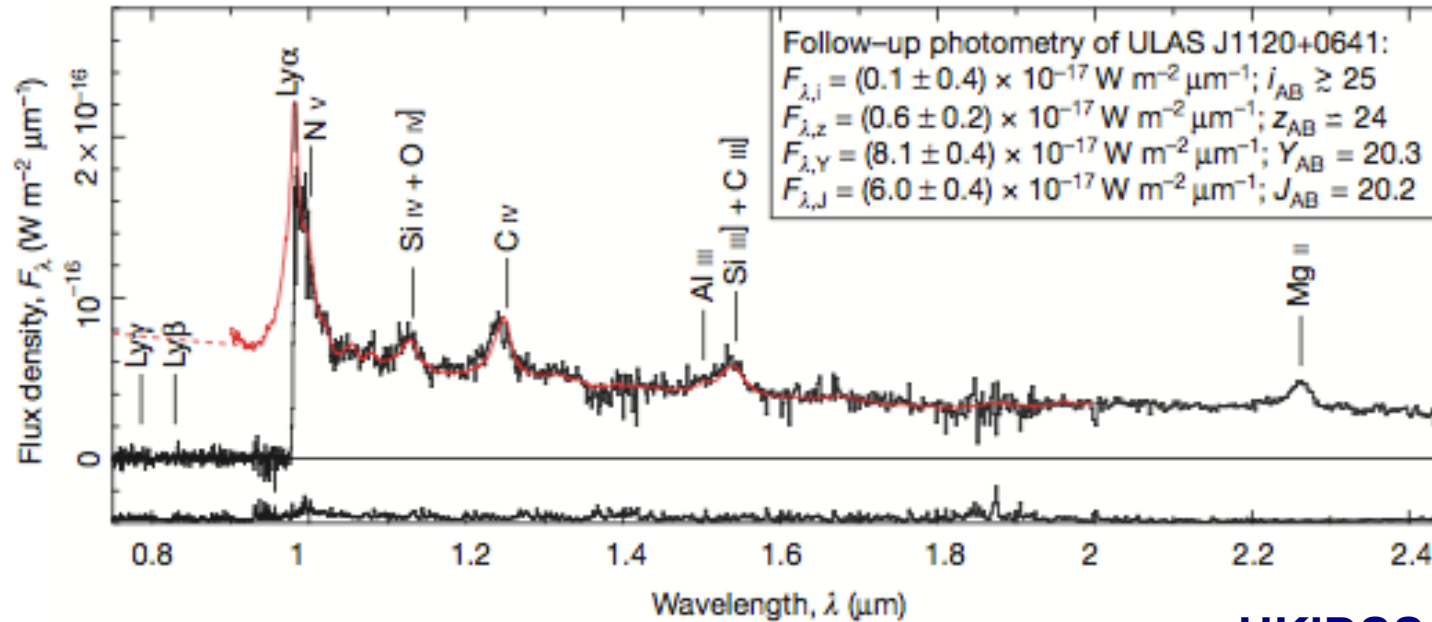
The first Black Holes

Black hole seeds

- Population III stars ($10^2 M_{\odot}$)
- Direct collapse of gas clouds ($10^{4-5} M_{\odot}$)
- Stellar mergers in dense clusters ($10^{4-5} M_{\odot}$)



ULASJ1120+0641, $z=7.08$



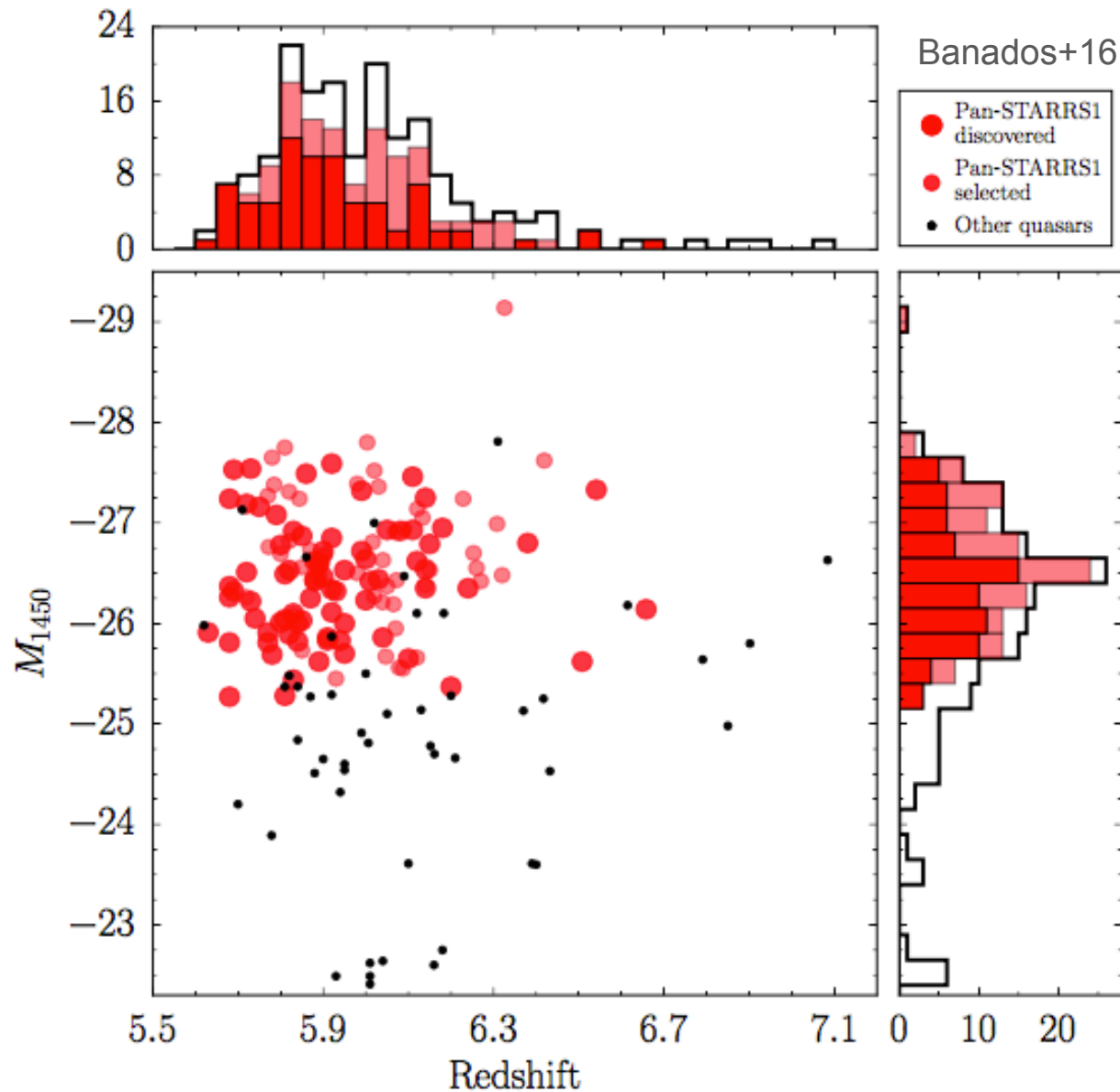
UKIDSS

Mortlock et al. 2011, GNIRS+FOR2, compared to average $z \sim 2.5$ SDSS QSOs
Fully mature QSOs at high redshift

See lessons on *The high-redshift Universe, and the role of galaxies and AGN to cosmic reionization* at

http://www.bo.astro.it/~vignali/PhD_Lessons/reionization_2015/

High-redshift quasars



Situation as of Sept. 2016

173 quasars known at
 $z > 5.6$

(≈ 80 at $z > 6$; 7 at $z > 6.5$)

172 optical-near-IR
selected

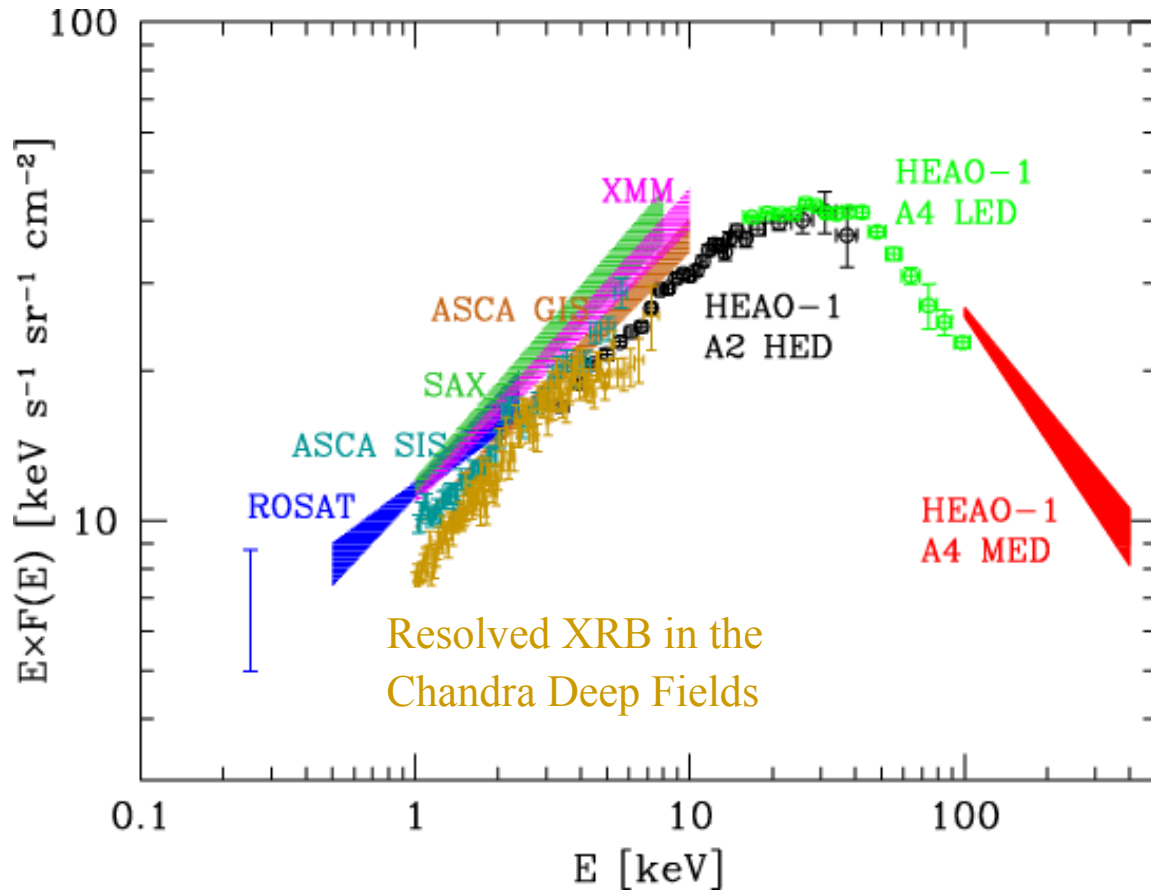
(Fan+00-06; Jiang+08-15;
Willott+07-10; Banados+14-16;
Mortlock+11; Venemans
+13-15; Matsuoka+16)

1 radio selected
(Zeimann+11)

0 X-ray selected

X-ray background, and the role of X-ray surveys

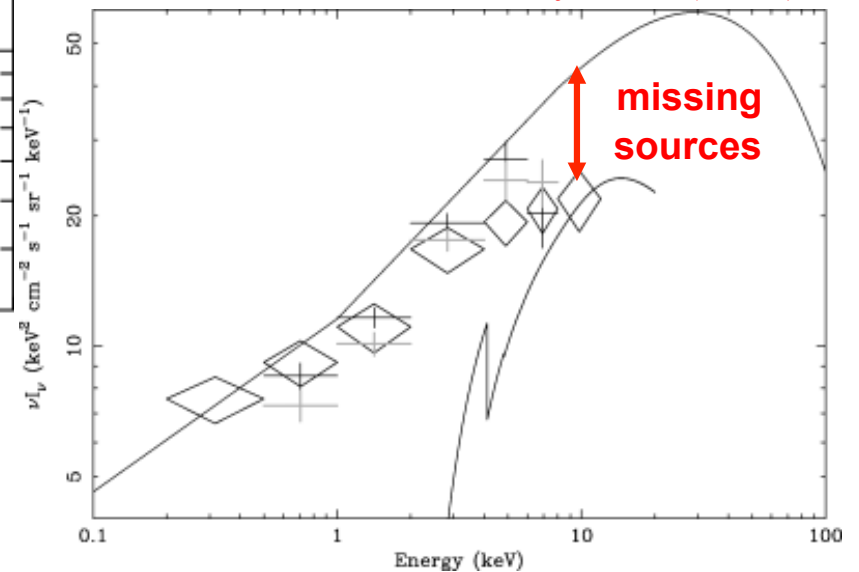
Resolved XRB fraction: still a “missing” population?



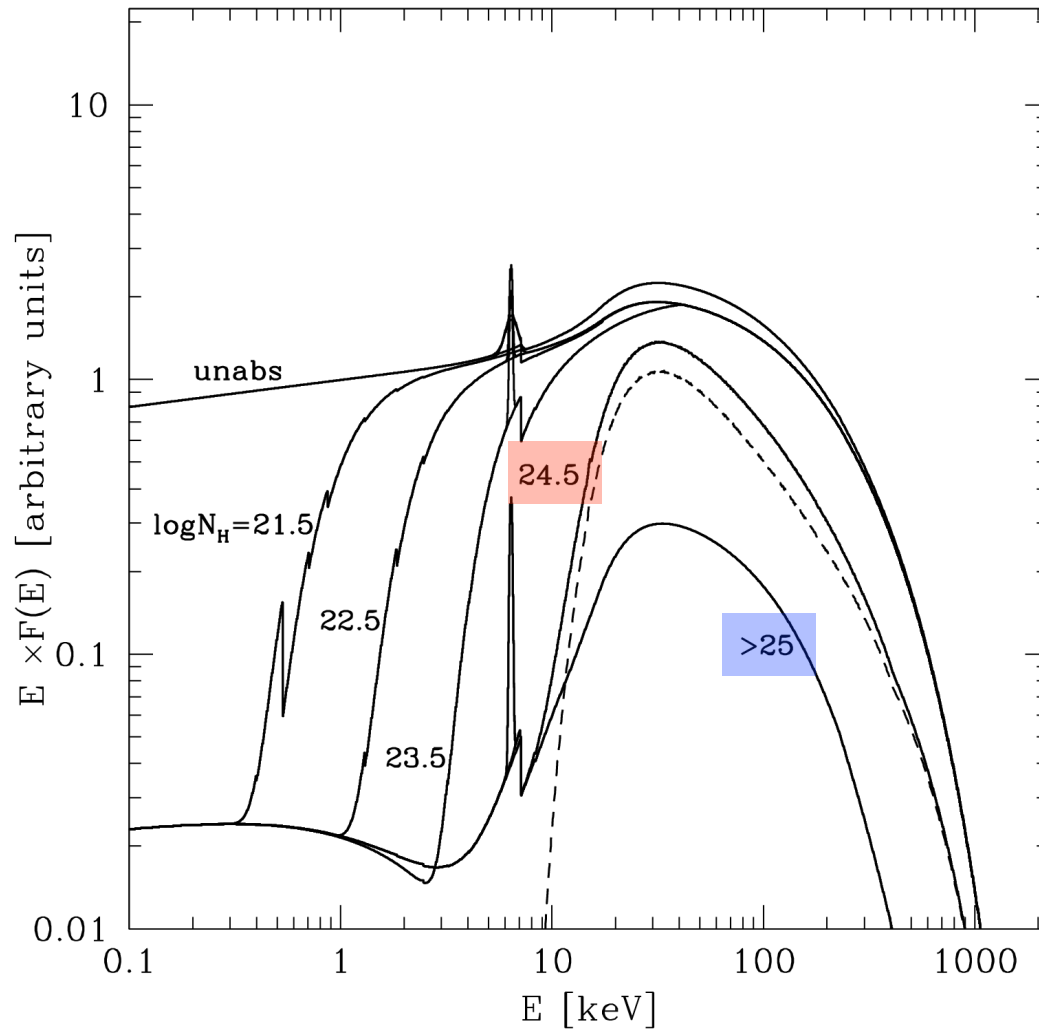
$\approx 80\text{-}95\%$ of the XRB being resolved into single sources at $E < \text{a few keV}$

BUT only 50% resolved above 5 keV

Worsley et al. (2005)



AGN X-ray spectral templates with different N_H

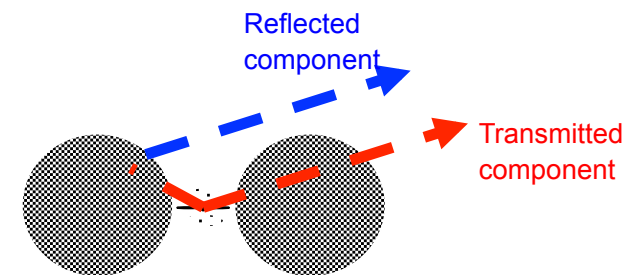


Only ≈ 40 -50 “secure” Compton-thick AGN (≈ 10 mildly-thick) known at present

Unabsorbed:
 $\log N_H < 21$

Compton-Thin
 $21 < \log N_H < 24$

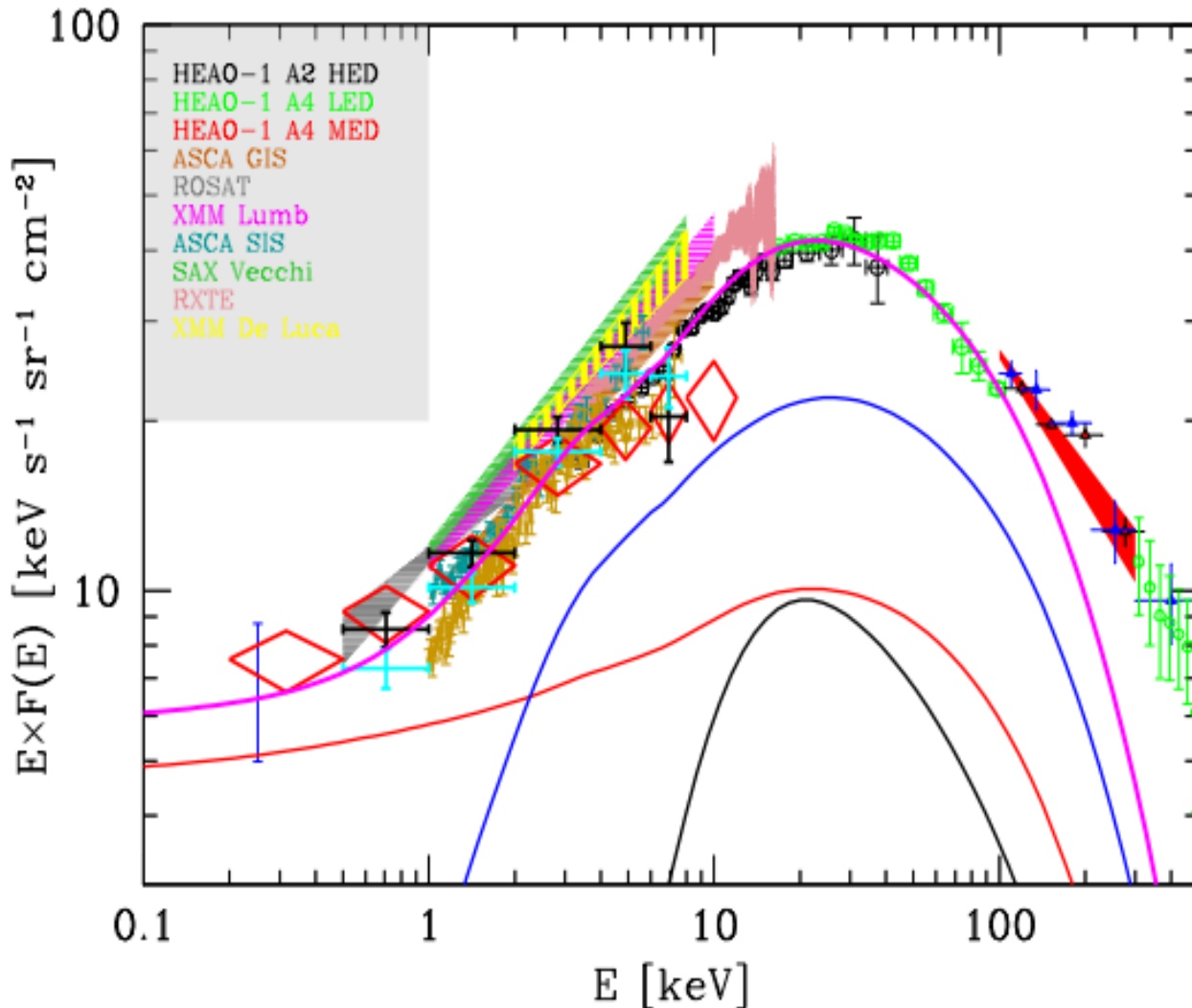
Compton-Thick:
Mildly ($\log N_H = 24$ -25)
Heavily ($\log N_H > 25$)



The cold gas in the torus contributes to the iron $K\alpha$ line emission.

As N_H increases, the spectrum is absorbed towards higher and higher energies.

The spectrum of the cosmic XRB as sum of obscured and unobscured AGN (following the original idea of Setti & Woltjer 1989)



The **XRB** synthesis provides an integral constraint (Gilli et al. 2007)

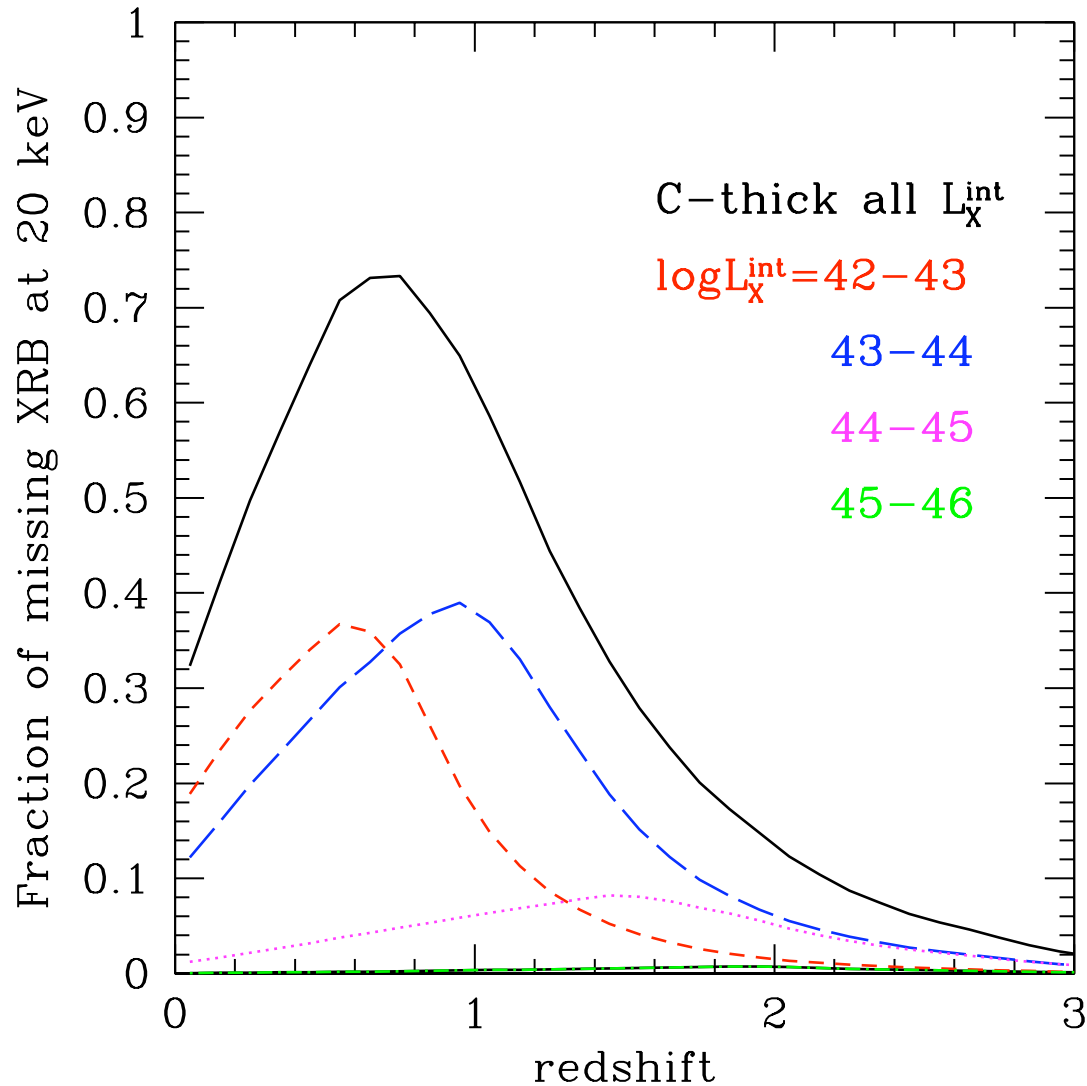
Red → unobscured

Blue → Compton Thin

Black → Compton Thick ($N_H > 10^{24}$ cm⁻²)

The evolution is folded in the adopted XLF

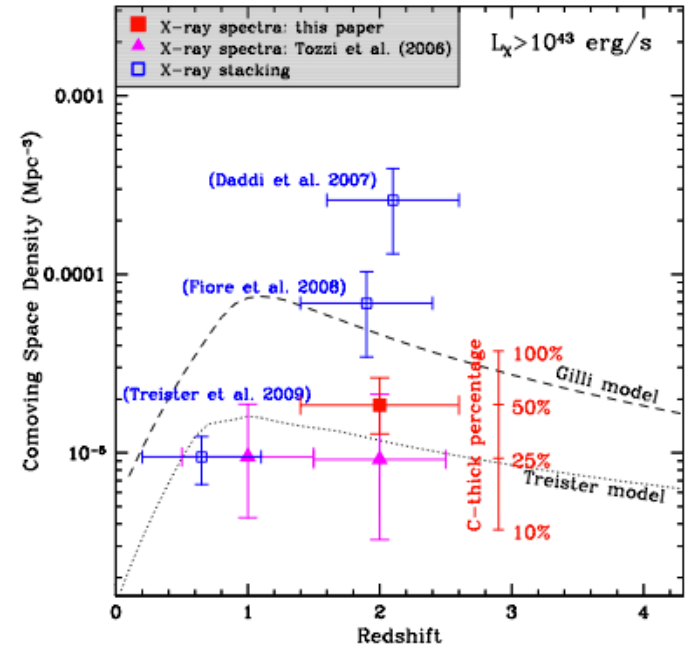
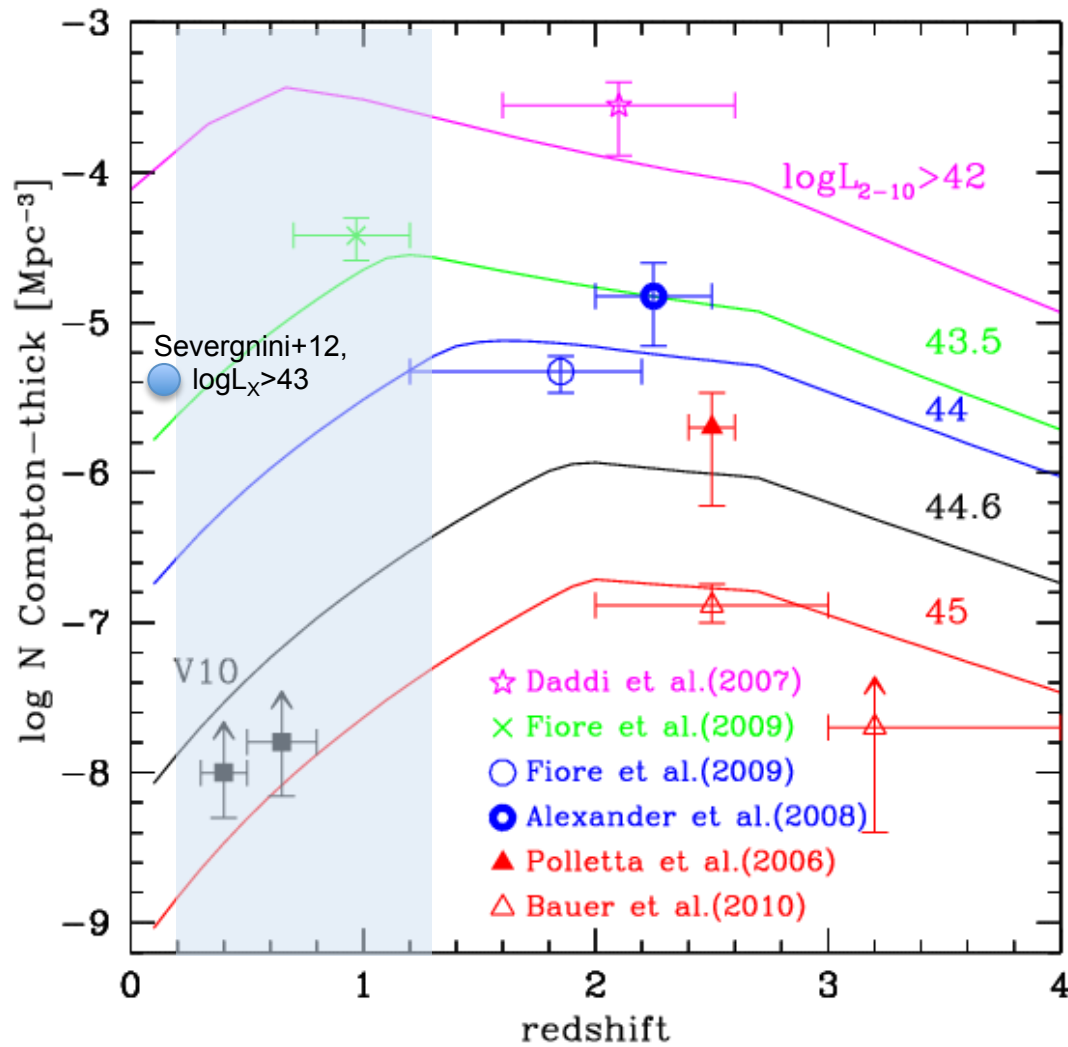
When the “missing” XRB was emitted?



Predicted peak
at $z \approx 0.7-0.8$
Mostly at
 $L_x \sim 10^{42-44}$ erg/s

Gilli 2013

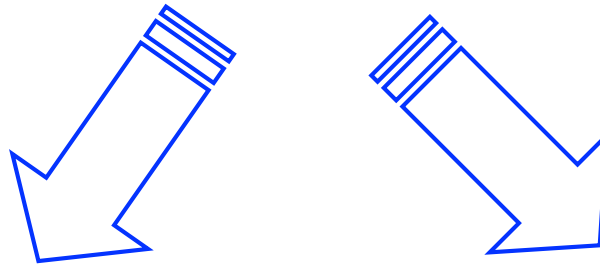
The space density of Compton-thick AGN



Alexander+11;
see also Bauer+10, ...

Vignali+10

Way to provide a census of AGN activity: X-ray surveys



Large-area survey

to pick up luminous and rare AGN

Relatively bright optical counterparts,
easier optical IDs

Deep-area survey

to pick up faint and distant AGN

Typically faint optical
counterparts, difficult optical IDs

What is the best observing strategy for X-ray surveys?

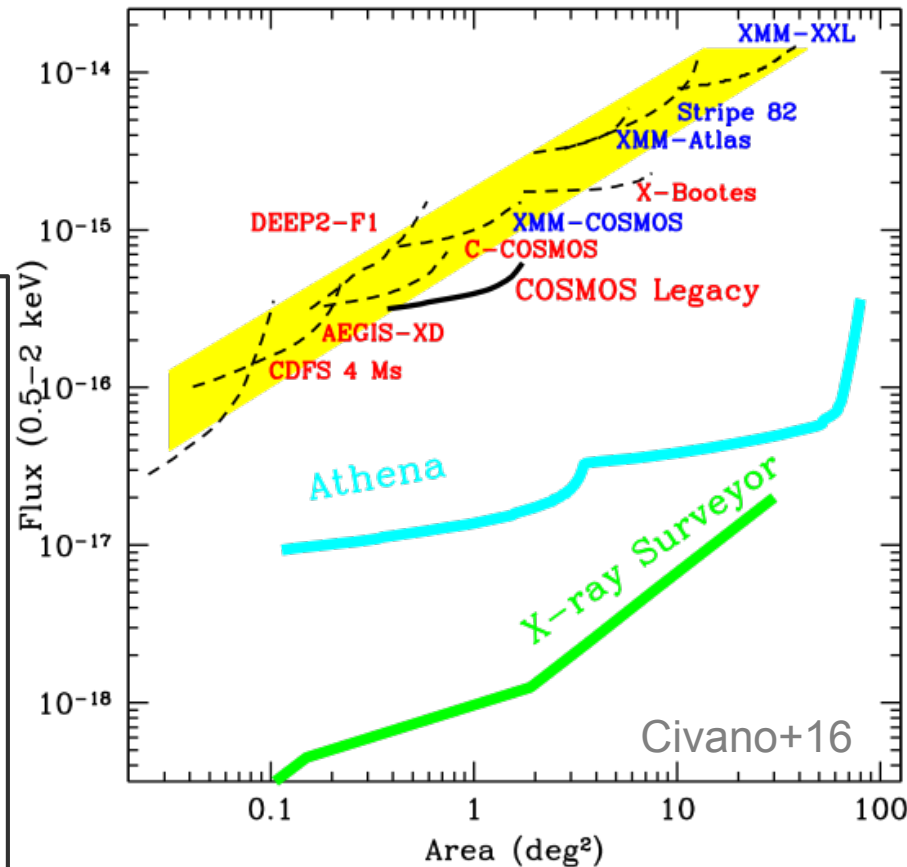
DEEP X-RAY SURVEYS

PROs:

- Ideal to reveal distant sources (because of the depth of the exposure)
- Large number of sources

CONs

- Limited to small areas
- Limited individual photon statistics



LARGE (and SHALLOW) X-RAY SURVEYS

PROs:

- Ideal to pick up bright and rare X-ray sources
- Possibility to cover large areas of the sky

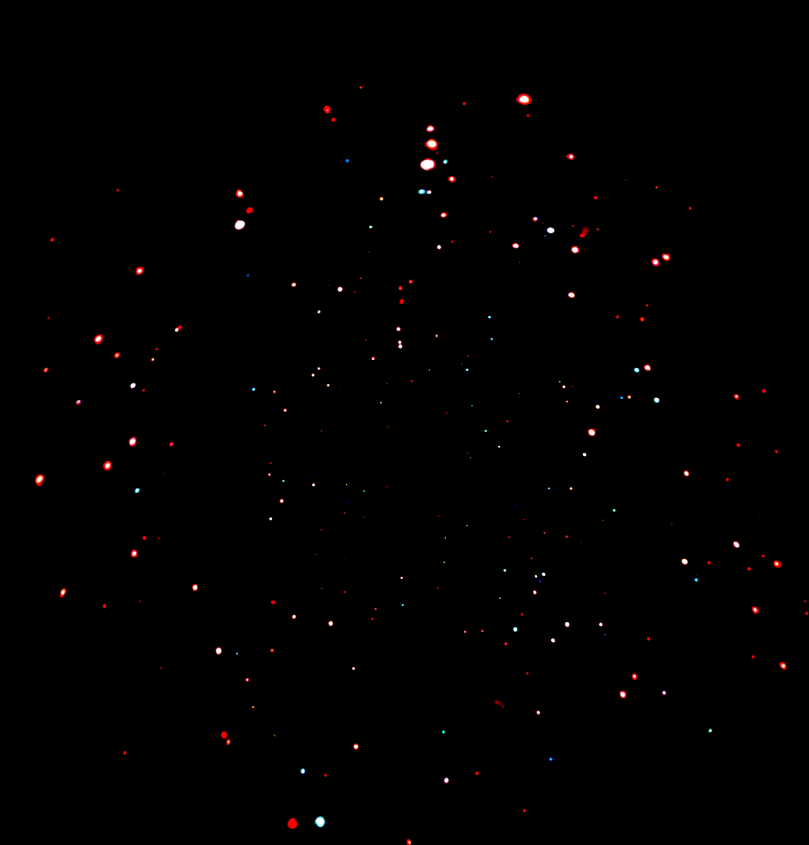
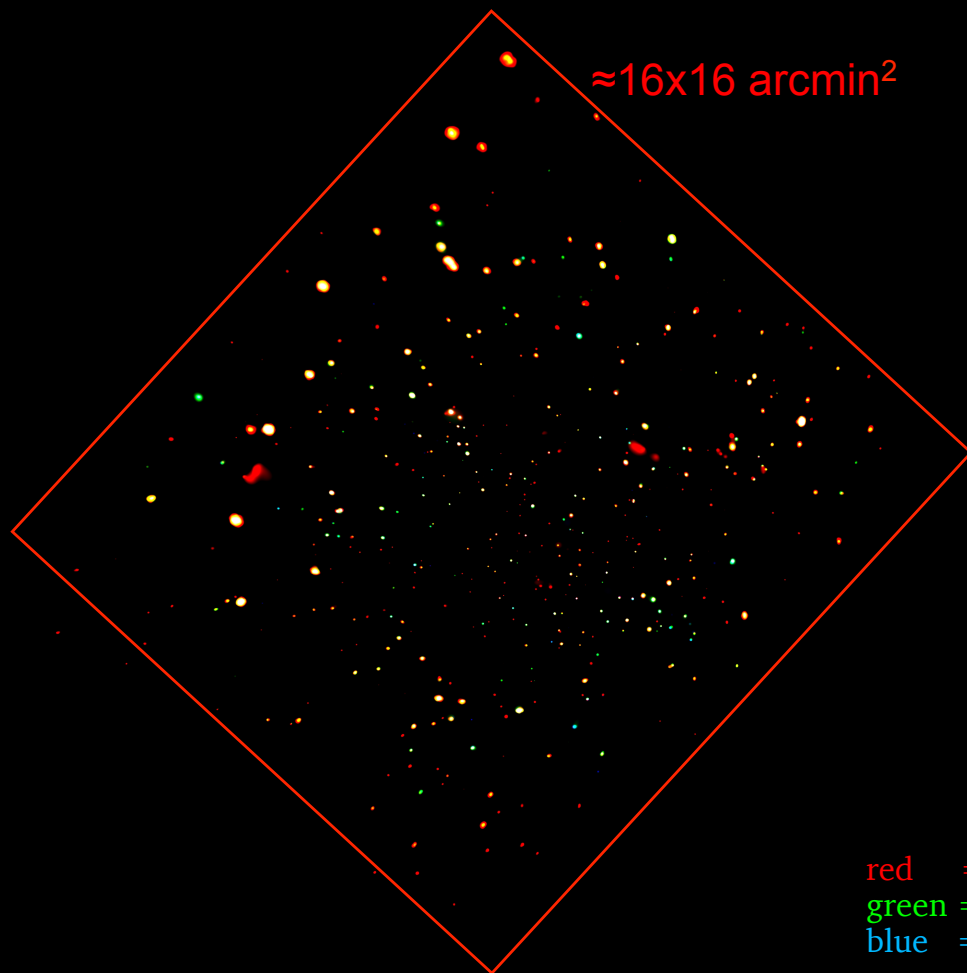
CONs

- Limited number of sources

Chandra Deep Fields

CDFN (Alexander+03; Luo+ 08,10)

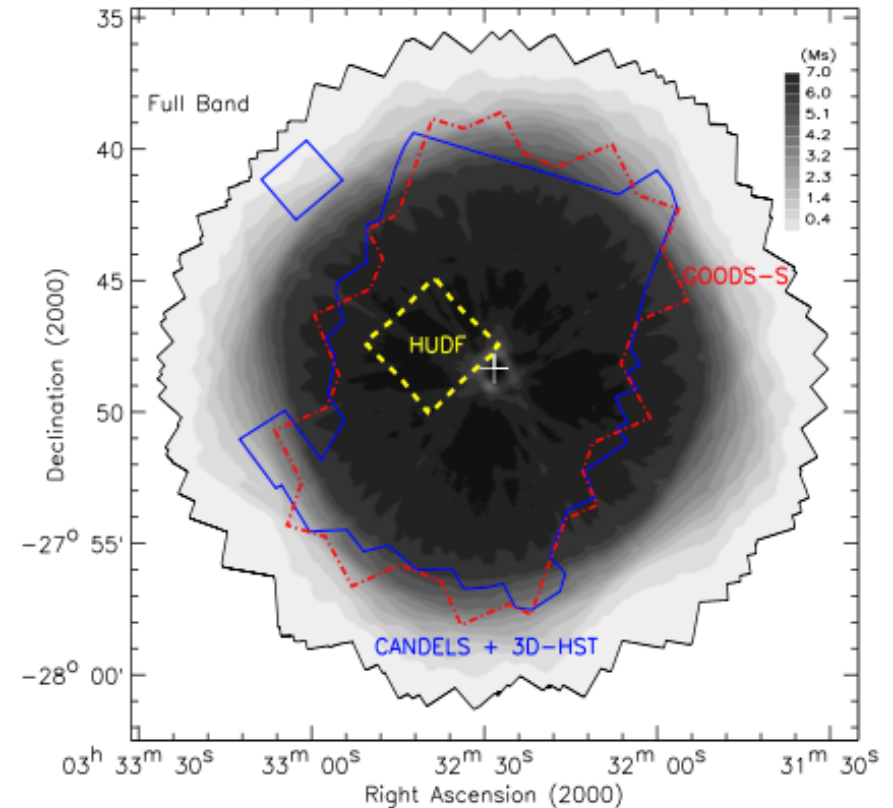
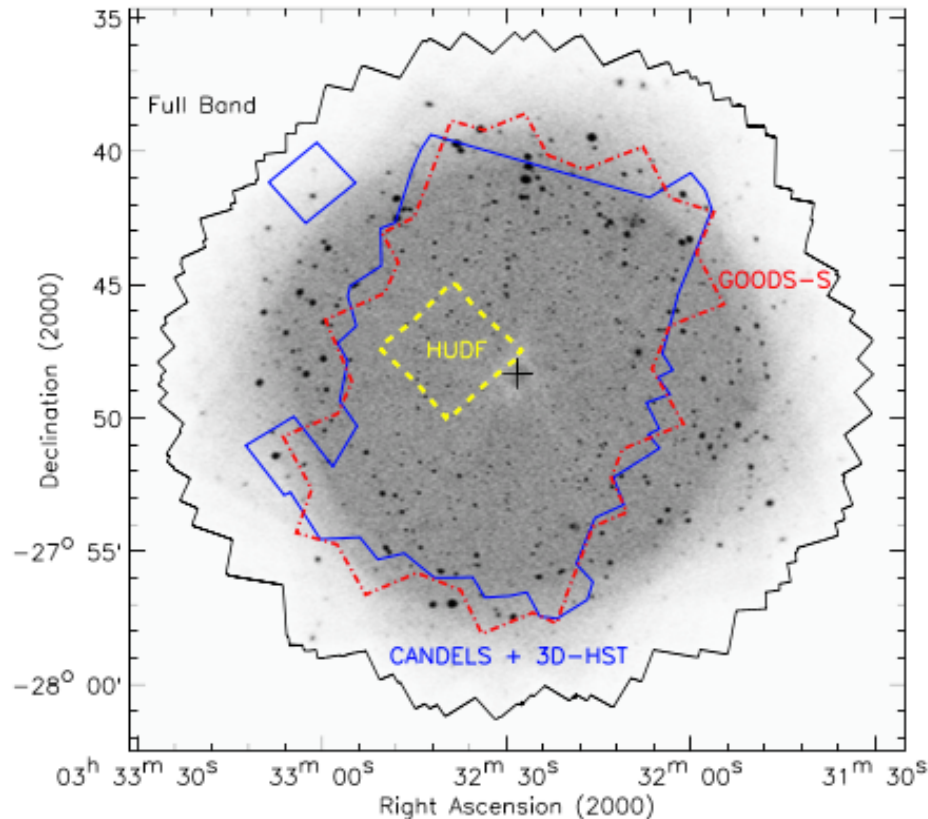
CDFS (Giacconi+02)



red = 0.5-1 keV
green = 1 - 2 keV
blue = 2 - 8 keV

The 7Ms *Chandra* Deep Field South. I.

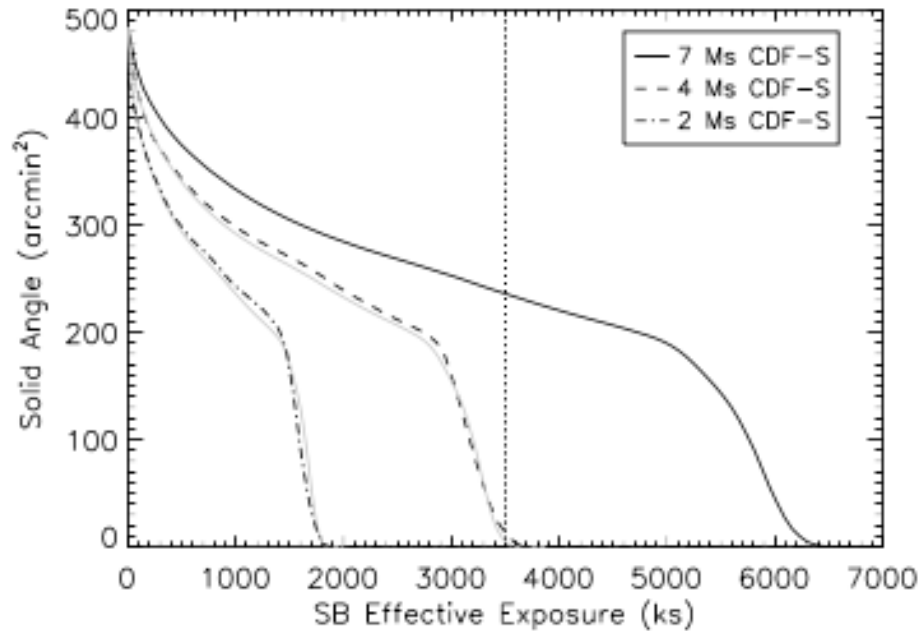
The deepest X-ray exposure ever



- 484 arcmin²
- 1008 X-ray sources (992 with counterpart, ≈66% with spec. redshift)
- At least 70% are classified as AGN
- Inner 1 arcmin region: $F_{[0.5-7\text{keV}]} = 1.9 \times 10^{-17}$ erg/cm²/s
 $F_{[0.5-2\text{keV}]} = 6.4 \times 10^{-18}$ erg/cm²/s
 $F_{[2-7\text{keV}]} = 2.7 \times 10^{-17}$ erg/cm²/s

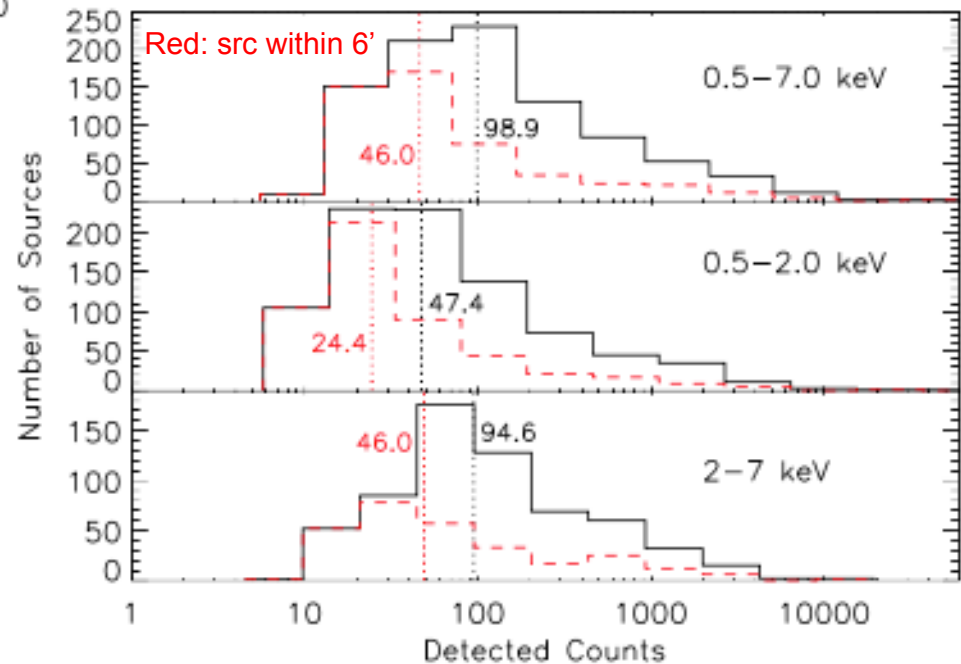
Luo+17

The 7Ms *Chandra* Deep Field South. II.



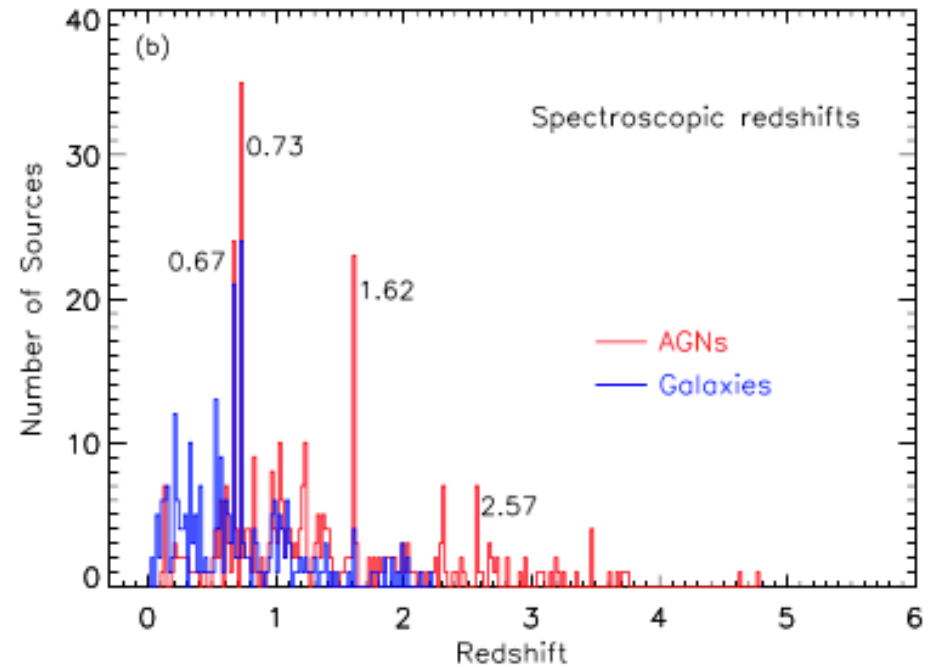
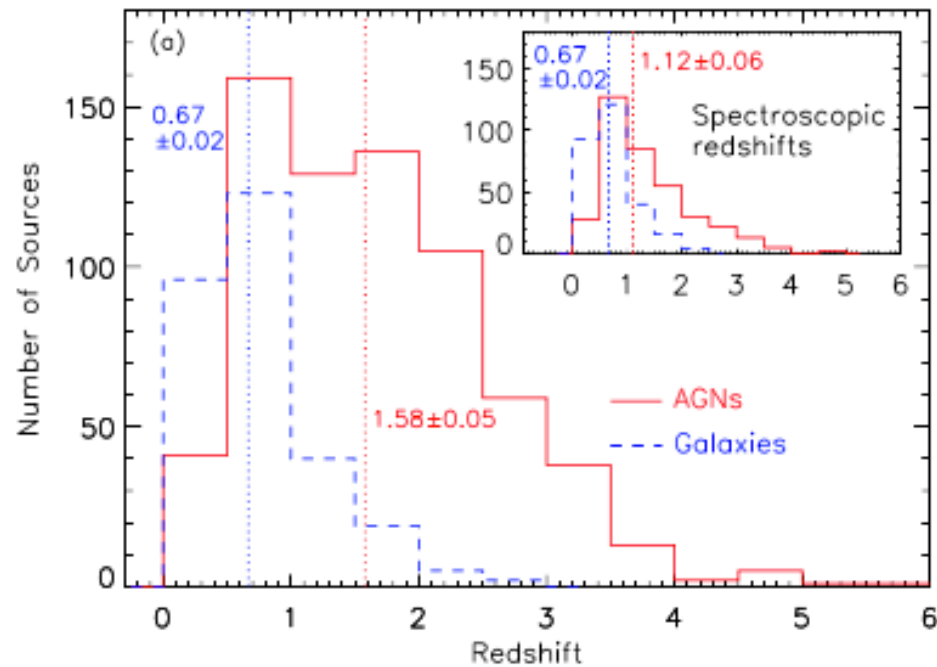
Solid angle vs. exposure time
Motivations behind going deeper

Number of counts
Median values around 100 (still low)



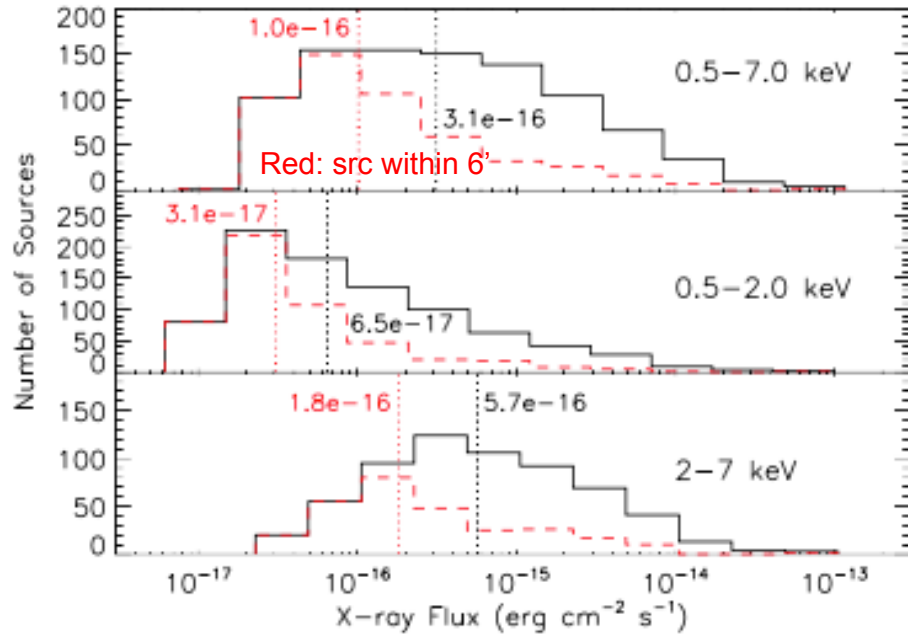
The 7Ms *Chandra* Deep Field South. III.

Redshift distribution AGN vs. Galaxies



Luo+17

The 7Ms *Chandra* Deep Field South. IV.

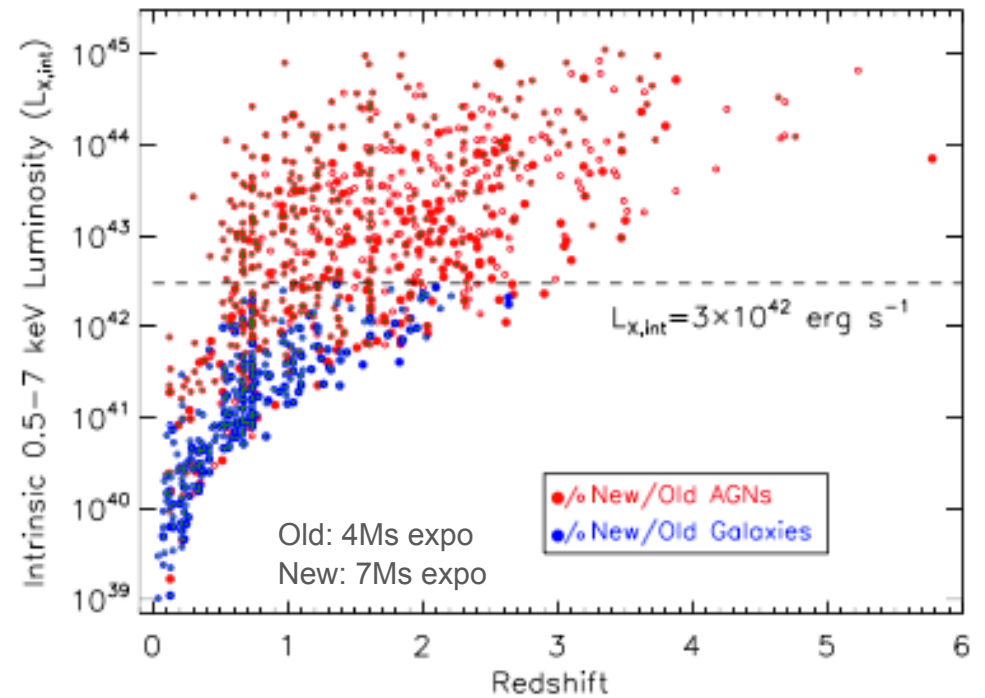


Flux distribution

Inner region vs. whole field

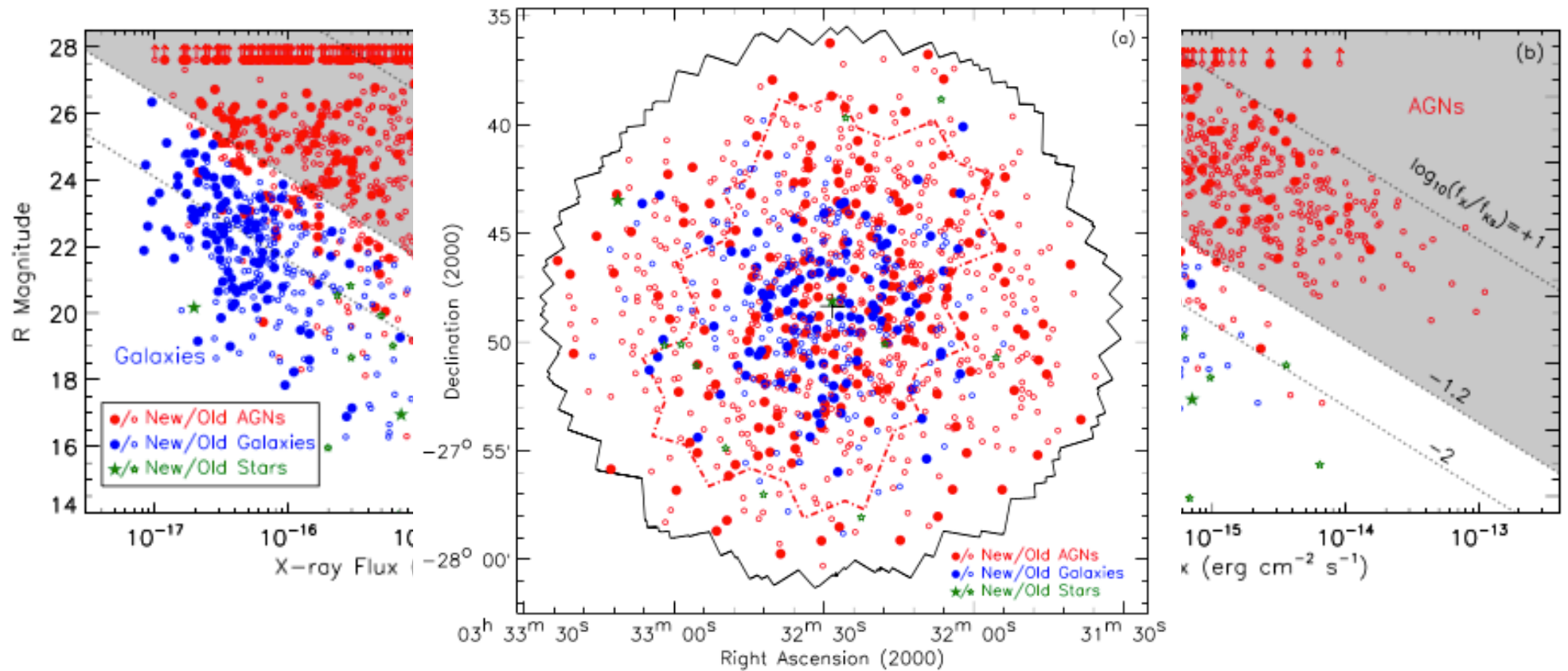
Luminosity distribution

AGN vs. Galaxies



The 7Ms *Chandra* Deep Field South. V.

R- (left panel) and K_S -band (right panel) mag vs. X-ray Flux



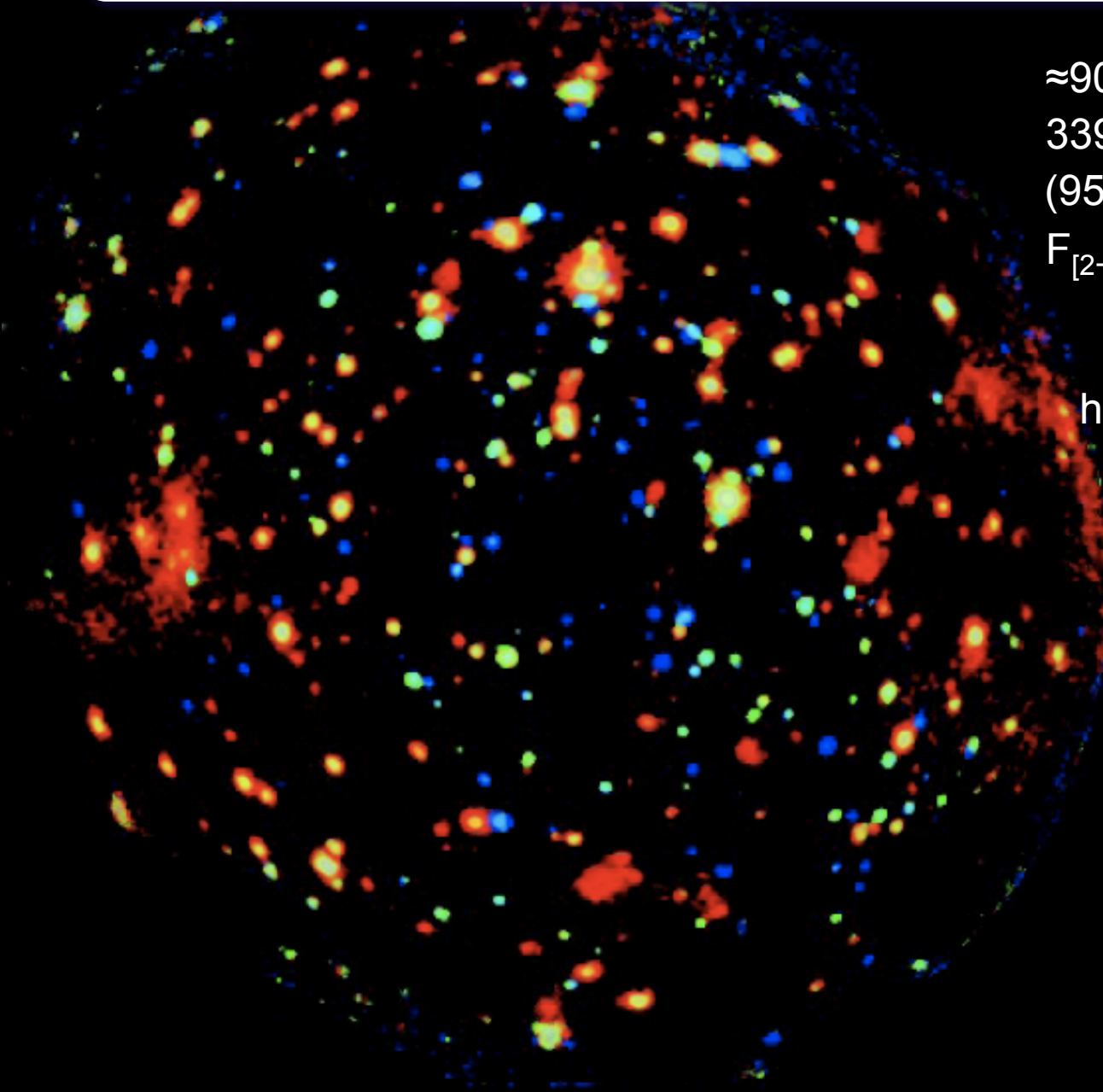
Chandra Deep Field South: XMM 3 Ms exposure

≈900 arcmin²

339 hard (2–10 keV) sources
(95% with spec/photo-z)

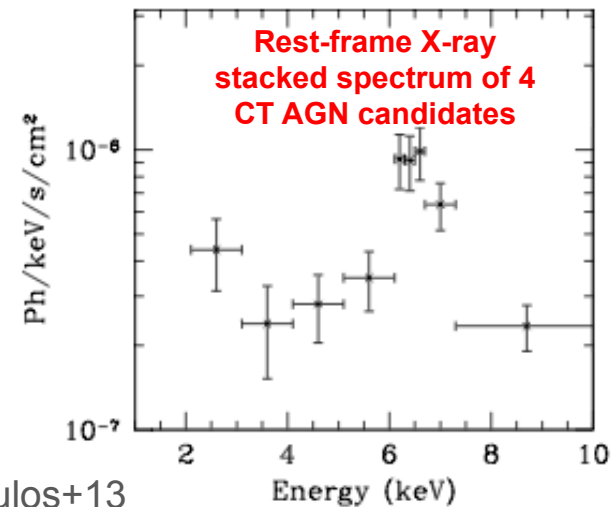
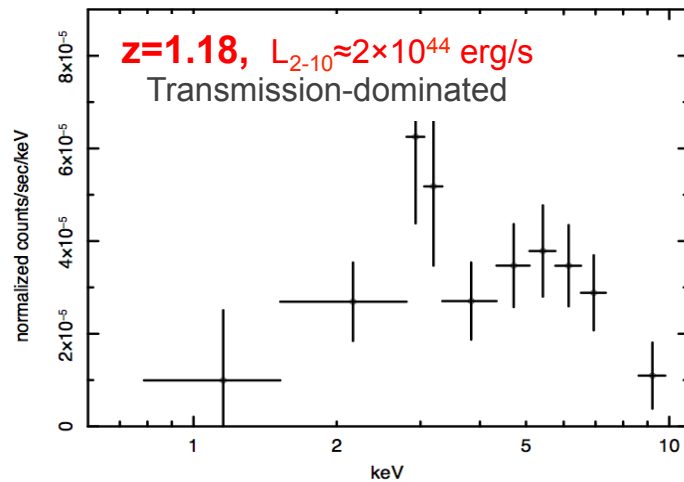
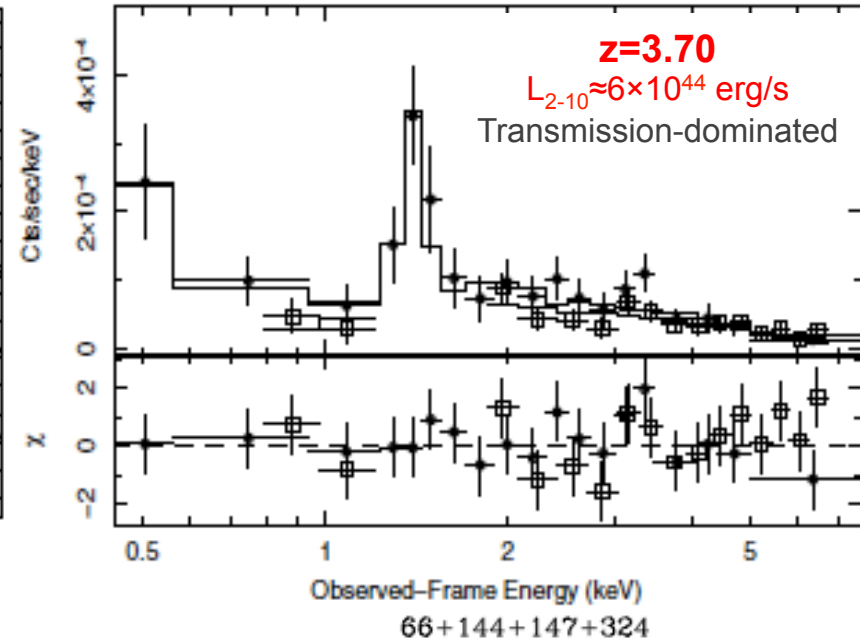
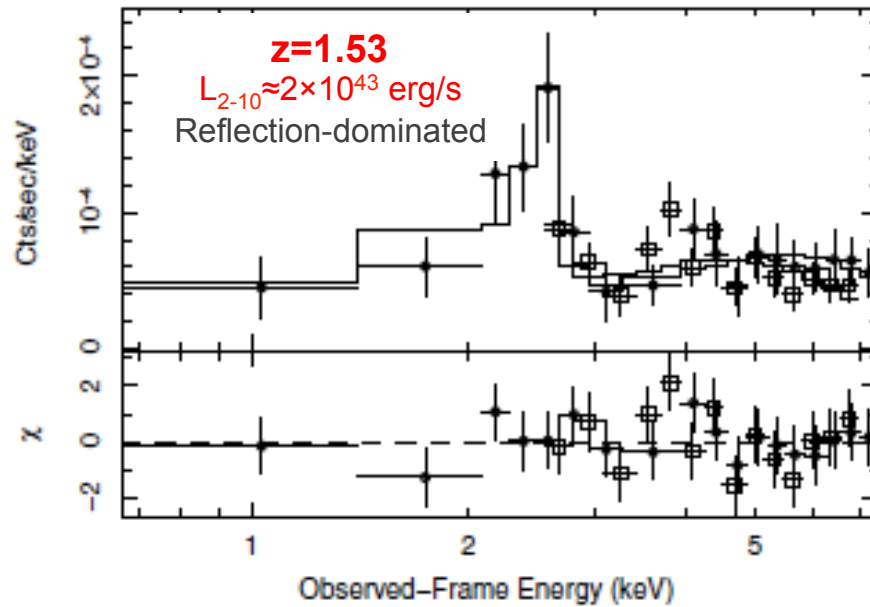
$F_{[2-10\text{keV}]} = 6.6 \times 10^{-16} \text{ erg/cm}^2/\text{s}$

Capable of probing the
high-z Universe with good
photon statistics



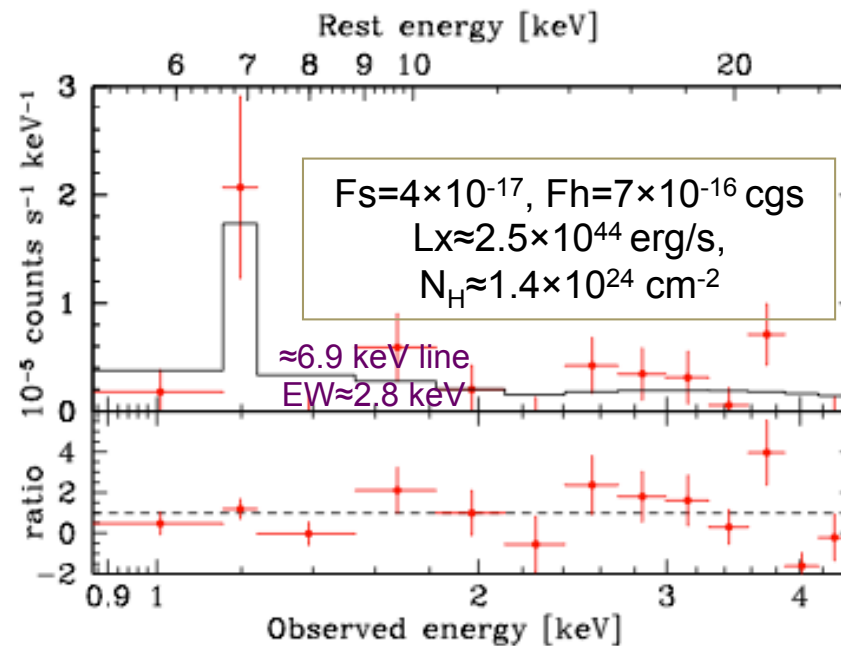
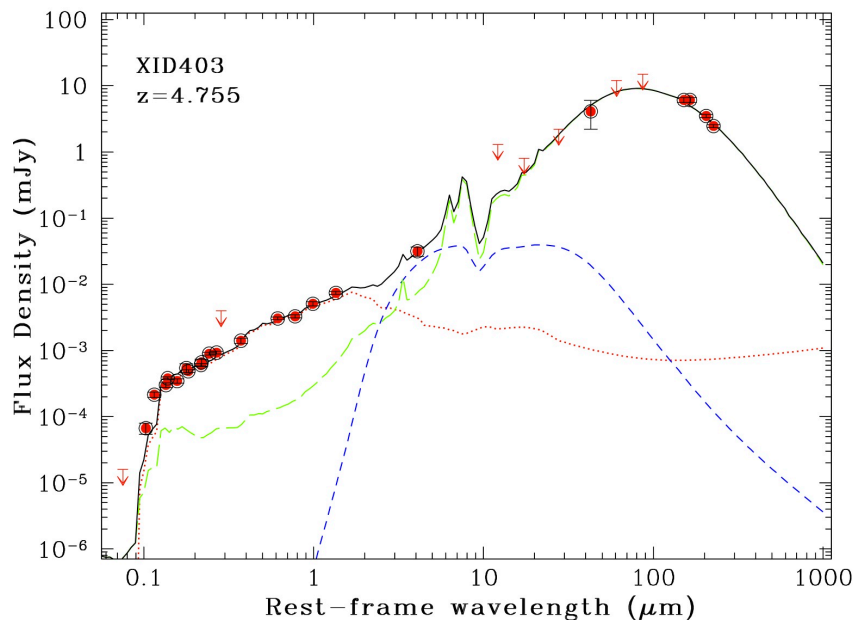
Distant obscured AGN in the CDF-S

Comastri+11



Georgantopoulos+13

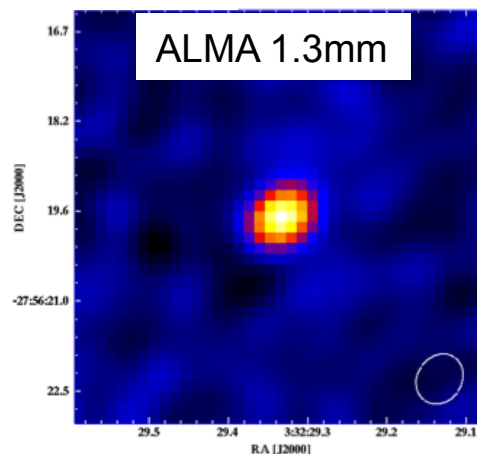
Obscured accretion and powerful star formation at $z=4.8$



$\text{SFR} \approx 1000 M_{\odot}/\text{yr}$
 $\Sigma_{\text{SFR}} > 26 M_{\odot}/\text{yr}/\text{kpc}^2$

Compact starburst, possibly
responsible for the X-ray
obscuration

Progenitor of compact quiescent
massive galaxies at $z \approx 3$



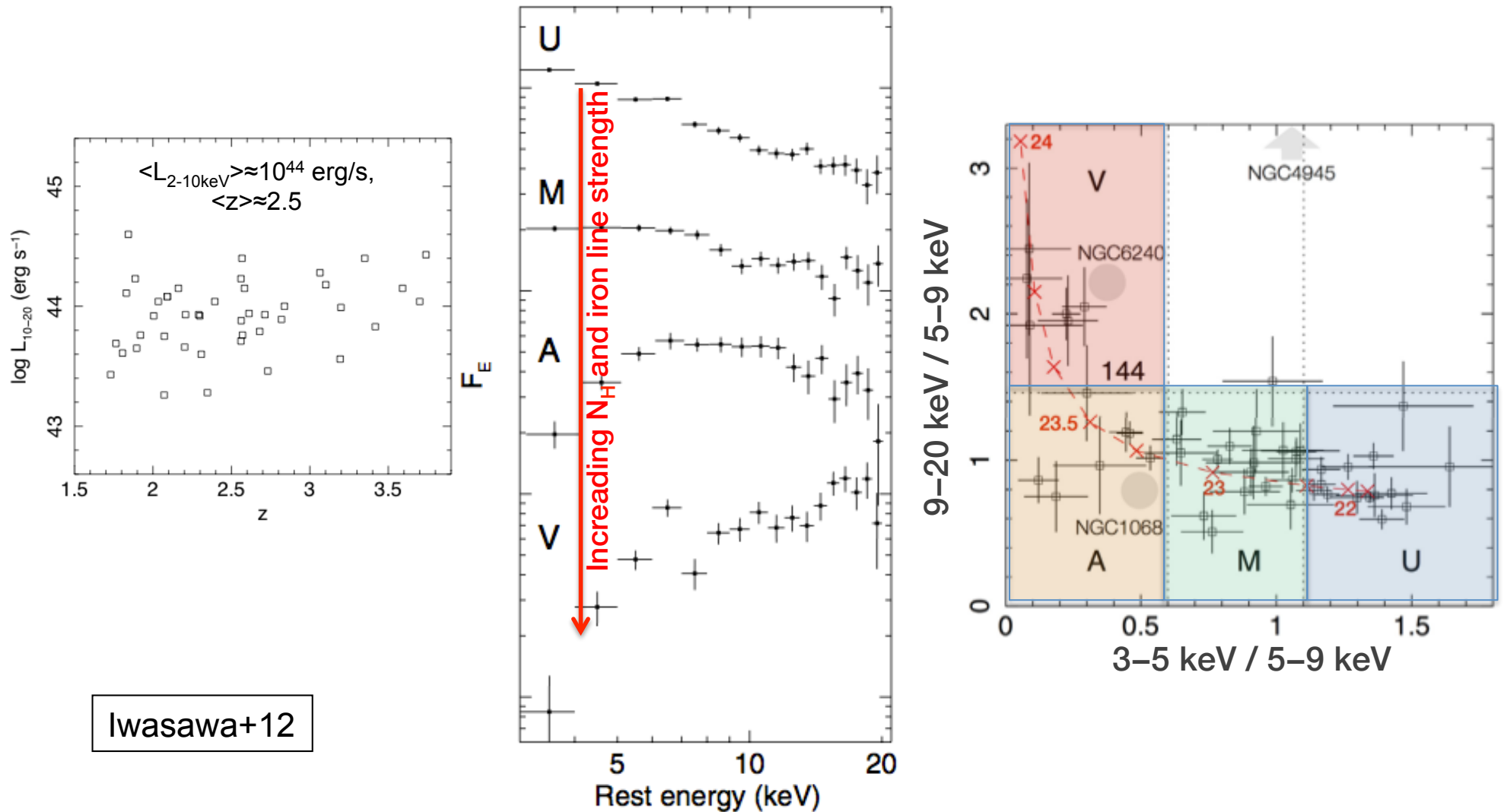
Gilli et al. 2014, 2011

Challenging and time-consuming
observations

This kind of studies are possible only
with deep ($> \text{Ms}$) *Chandra* exposures

X-ray stacking to probe high-redshift sources

At $z > 1.7$, the rest-frame 10–20 keV band enters the XMM-Newton bandpass → search for obscured AGN using the hard (>10 keV) excess – 44 AGN in the 2–10 keV source catalog



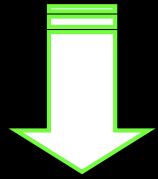
Iwasawa+12

X-ray variability from deep X-ray surveys

X-raying the COSMOS

XMM-Newton
1.55 Ms
1822 sources

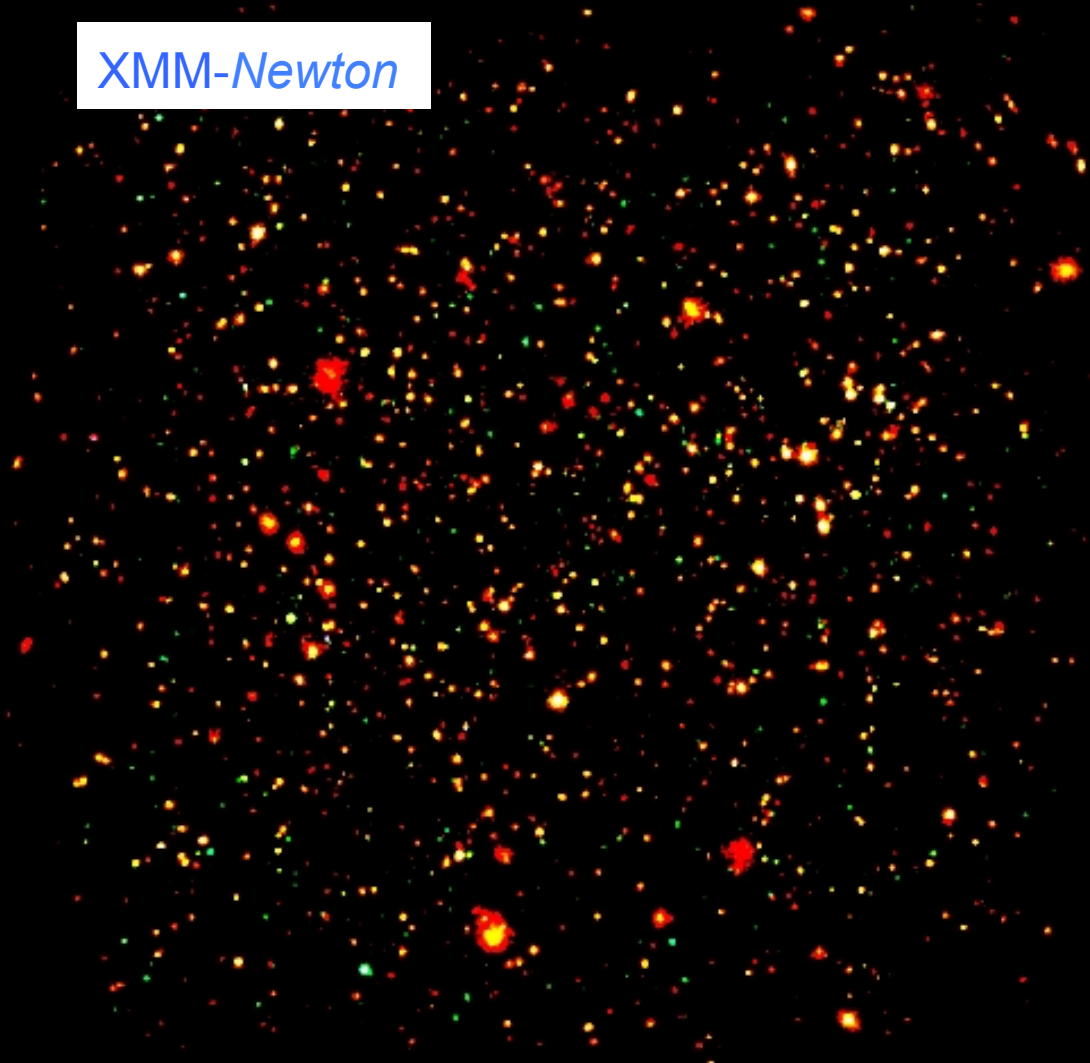
Large area, 2 deg², good
photon statistics
Background and PSF size
main limitations



Chandra
1.8 Ms
1761 sources

Guarantees depth on an area
initially of 0.9 deg²

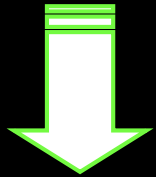
XMM-Newton



X-raying the COSMOS

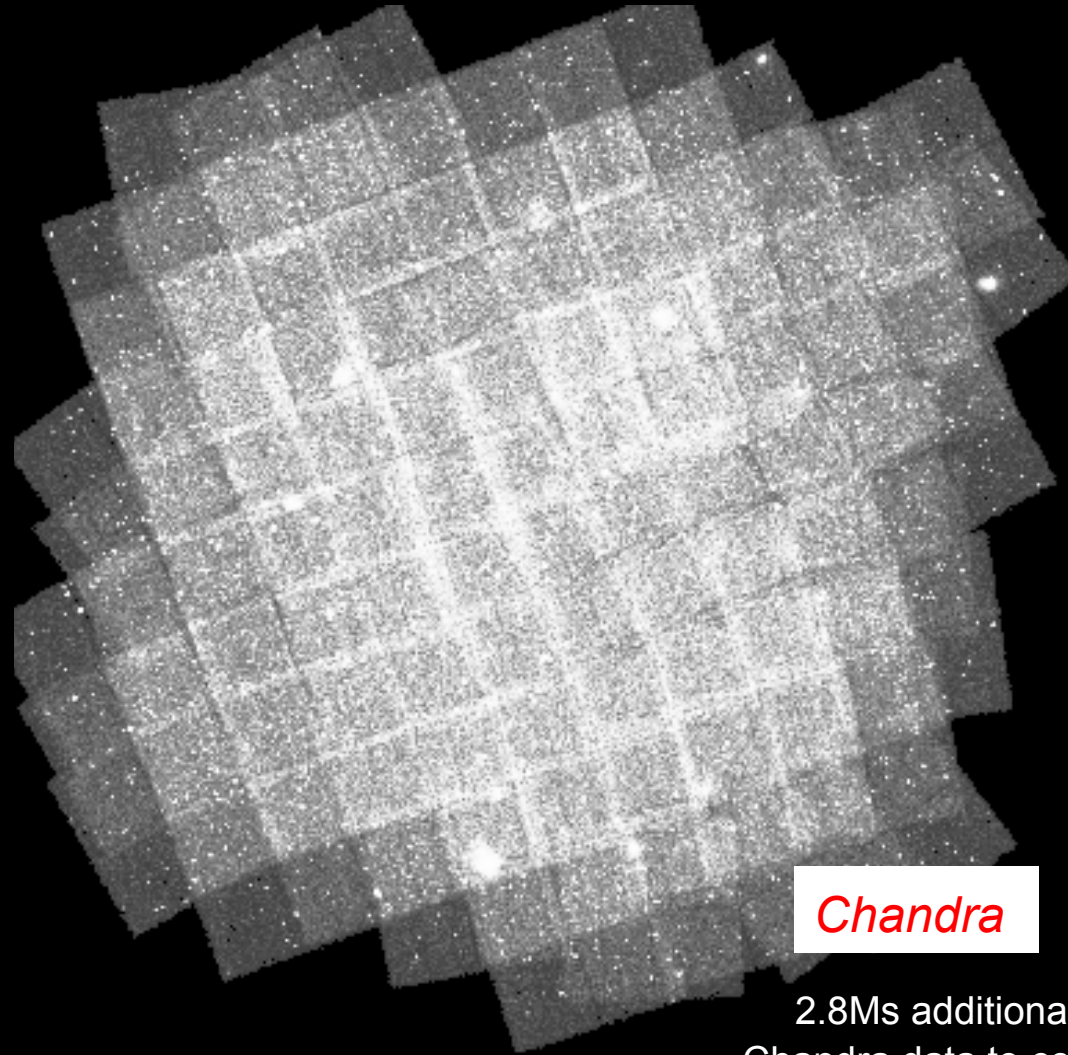
XMM-Newton
1.55 Ms
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Background and PSF size
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Chandra
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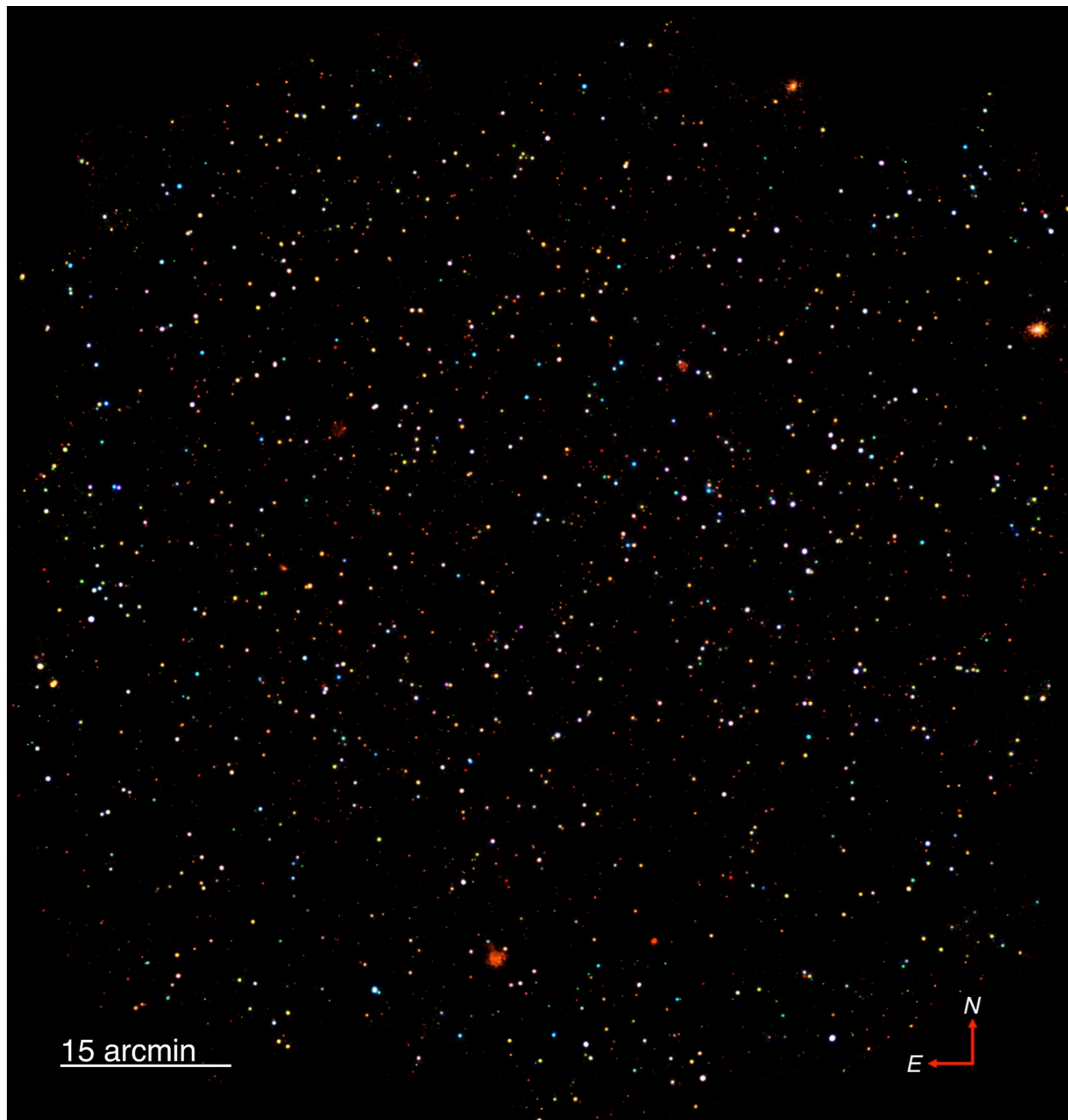
Chandra

2.8Ms additional
Chandra data to cover
the entire ~2.2 deg² (going deep)
4016 point-like X-ray sources

C
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Courtesy of
the Chandra
COSMOS
Legacy
collaboration

AGN Evolution

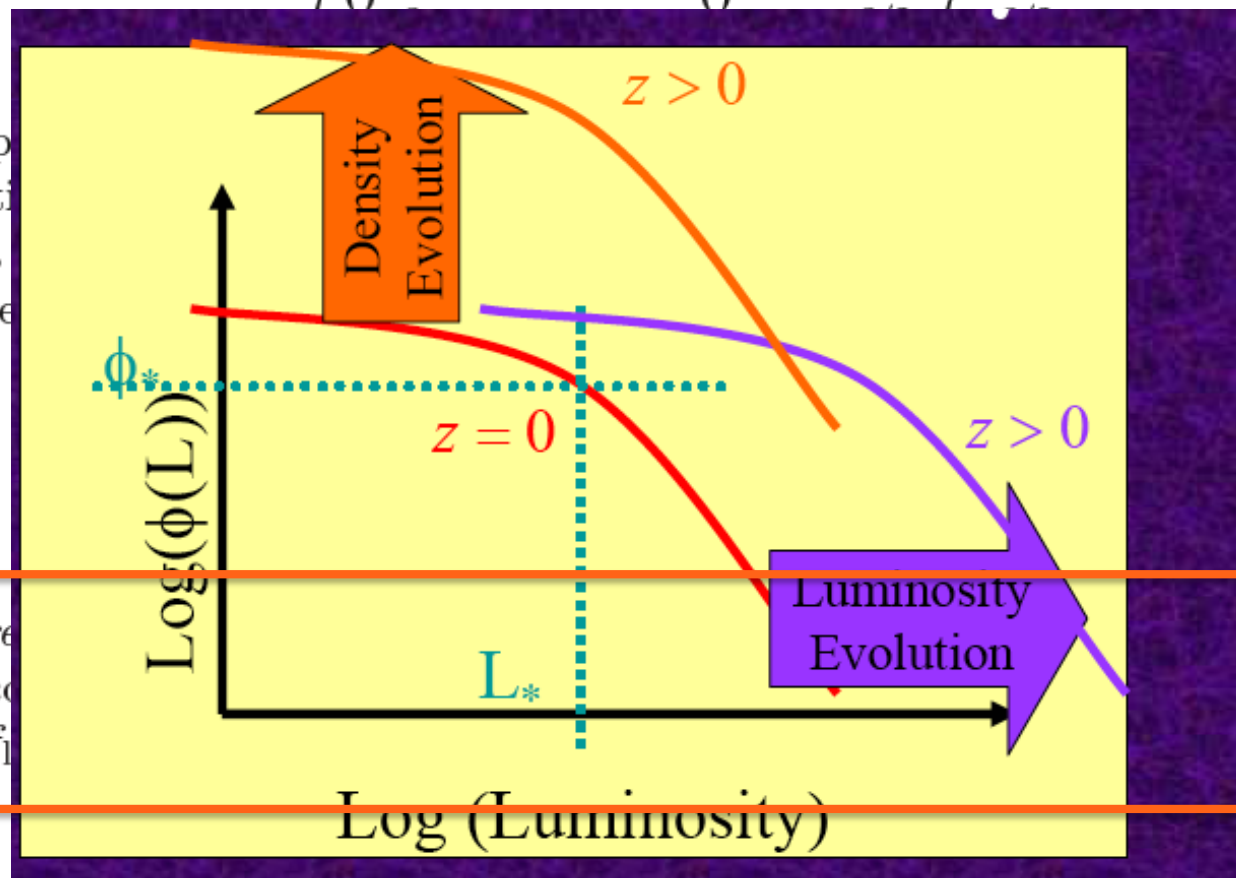
AGN surveys, basic definitions

The space density of sources of different intrinsic luminosities, L , is described by the *luminosity function* (LF), $\phi(L)$, so that $dN = \phi(L)dL$ is the number of sources per unit volume with luminosity in the range L to $L + dL$. Let us consider, for simplicity, the local or nearby (Euclidean) universe uniformly filled with sources with LF $\phi(L)$. If S is the limiting flux that we can detect, sources with luminosity L can be observed out to a distance $r = (L/4\pi S)^{1/2}$. The number of sources over the solid angle Ω , observable down to the flux S are:

$$\int_{L_{\min}}^{L_{\max}} \phi(L) \Omega r^2 dL \quad (1)$$

Thus, independent of a normalization of sources in a Ω volume, the number density of sources in a Ω volume varies as $L^{-3/2}$ (if we use $r = (L/4\pi S)^{1/2}$).

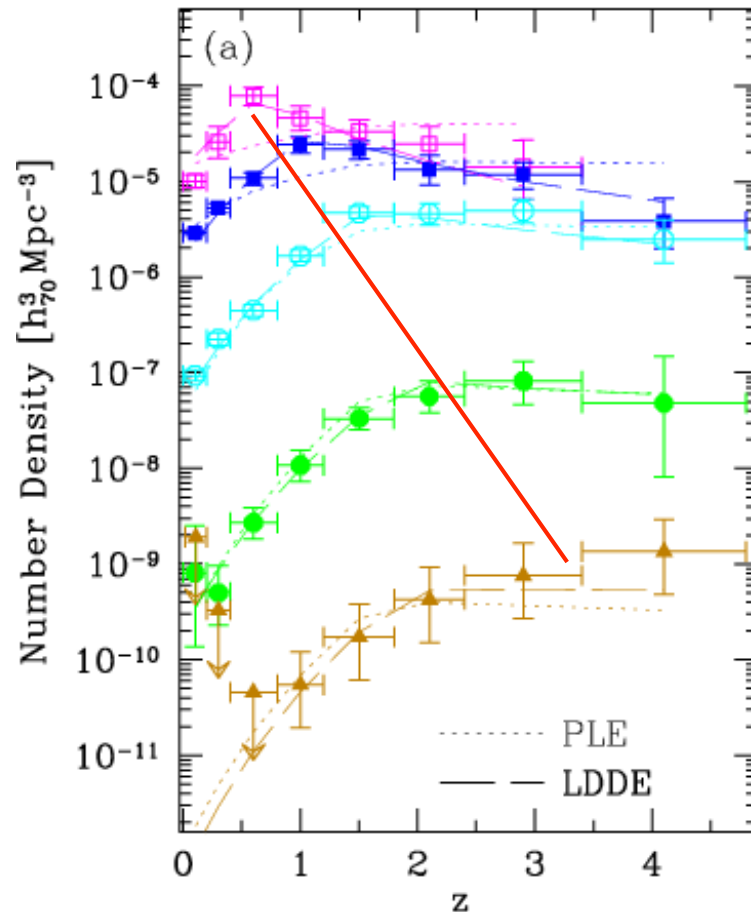
the determination of a non-evolving class of sources $N(> S)/d \log S = \int_{L_{\min}}^{L_{\max}} \phi(L) \Omega r^2 dL / d \log S = \Omega \int_{L_{\min}}^{L_{\max}} \phi(L) L^{-3/2} dL / d \log S = \Omega \int_{L_{\min}}^{L_{\max}} \phi(L) L^{-1/2} dL / d \log S = \Omega \int_{L_{\min}}^{L_{\max}} \phi(L) L^{-1/2} dL / (dL/L) = \Omega \int_{L_{\min}}^{L_{\max}} \phi(L) L^{1/2} dL$ (2)



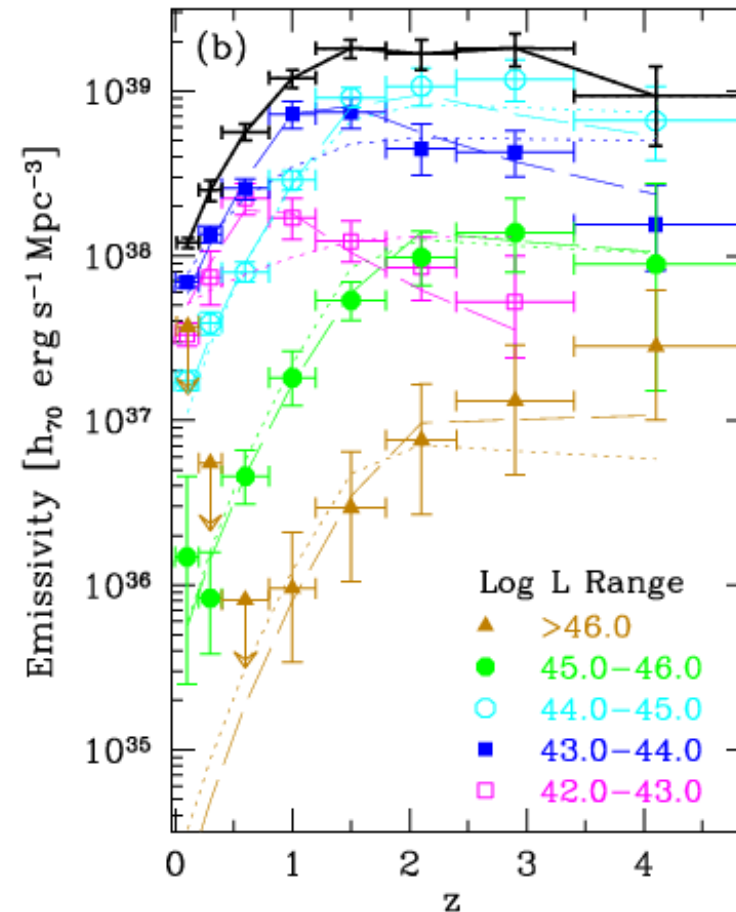
In the *pure* case of sources is constant (PDE) case (flux density evolution density evolution density varies.

AGN cosmological evolution

Number density



Luminosity density

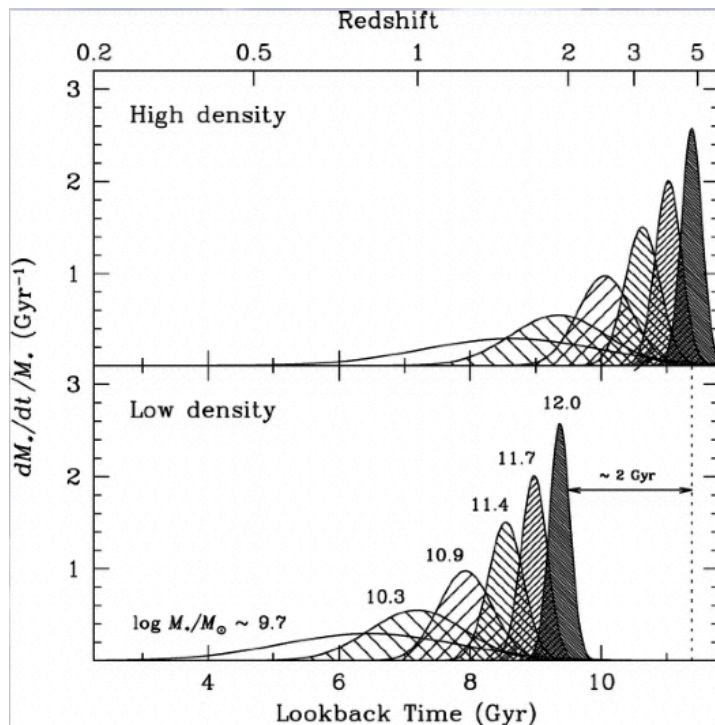


Objects with lower luminosity peak at lower redshift, similar to what observed for SFR in galaxies \Rightarrow **cosmic downsizing**
QSOs peak at $z \approx 2-3$, AGN at $z \approx 0.5-1$

The number density of AGN evolves differently for sources of varying luminosities
→ LDDE (luminosity-dependent density evolution) is the current, widely accepted parameterization of AGN evolution in X-rays

The density of the most luminous AGN peaks earlier in cosmic time than for less luminous objects, which likely implies that large black holes are formed earlier than their low-mass counterparts

Similar behavior for galaxies: massive galaxies tend to form stars earlier and faster than less massive galaxies



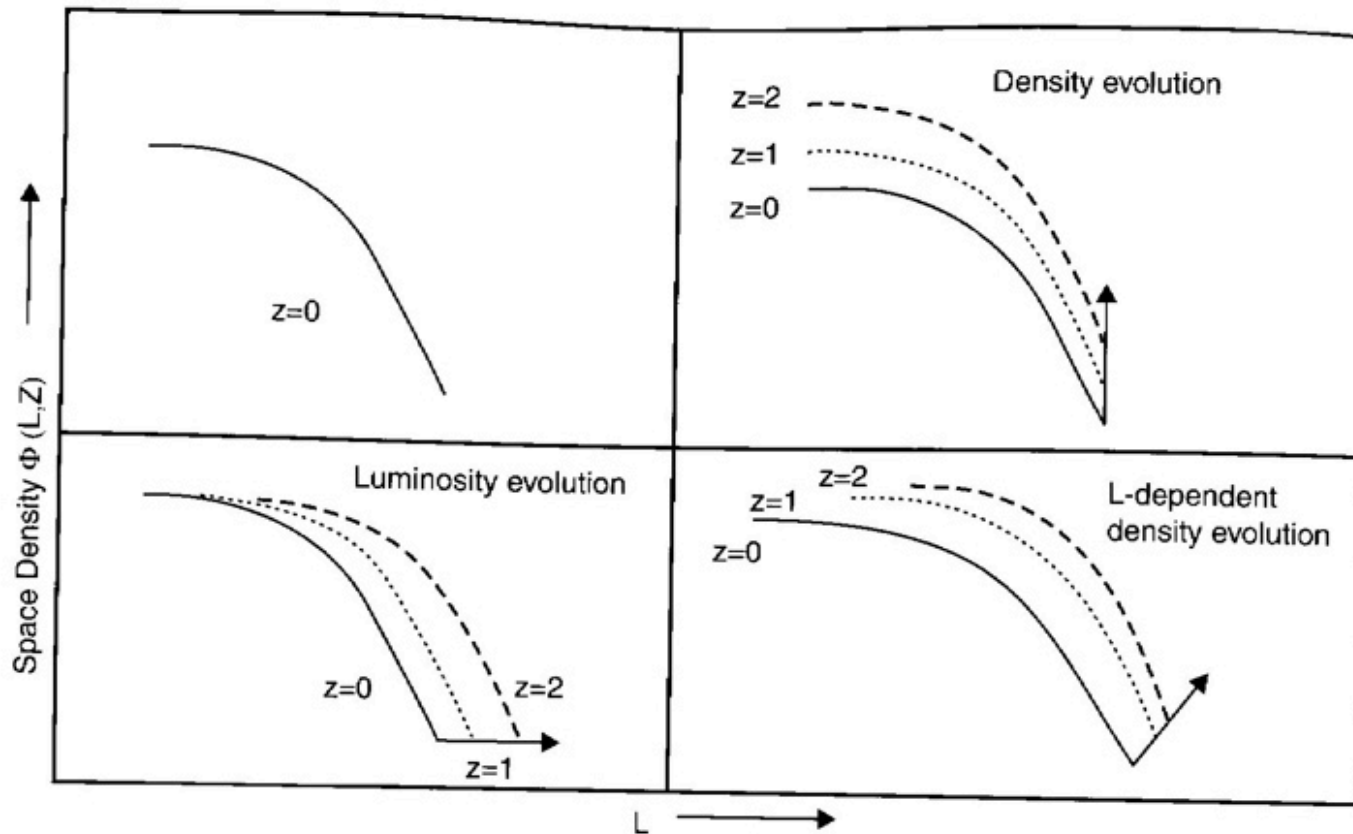
Galaxy formation took place in “downsizing”, with more massive galaxies forming at higher redshift (Cowie+96)

AGN and galaxies seem to share a similar behavior in terms of evolution

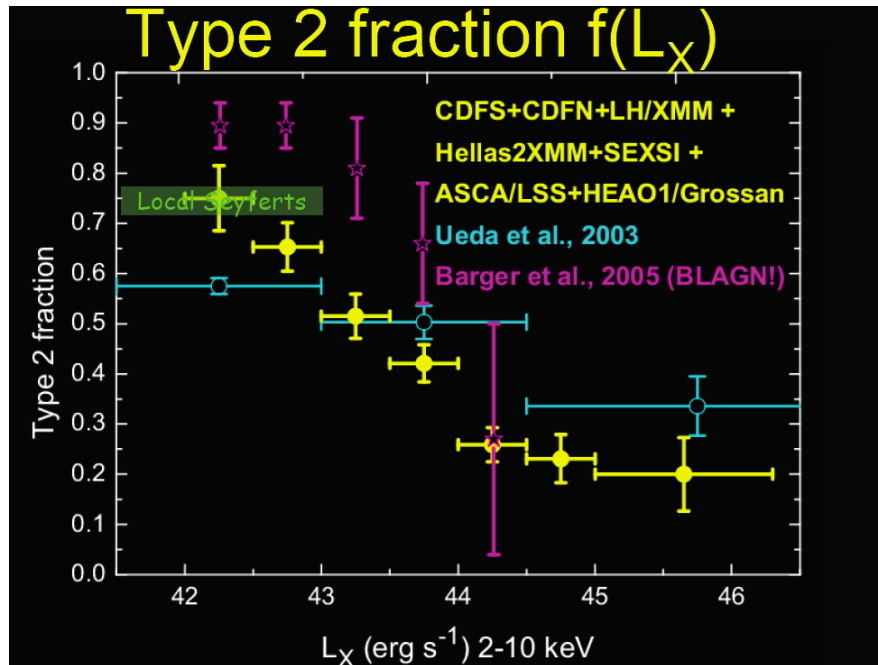
Luminosity Evolution:
AGN more luminous in the past

Density Evolution:
AGN more numerous in the past

Luminosity-dependent Density
Evolution:
Evolution in density dependent on
AGN luminosity

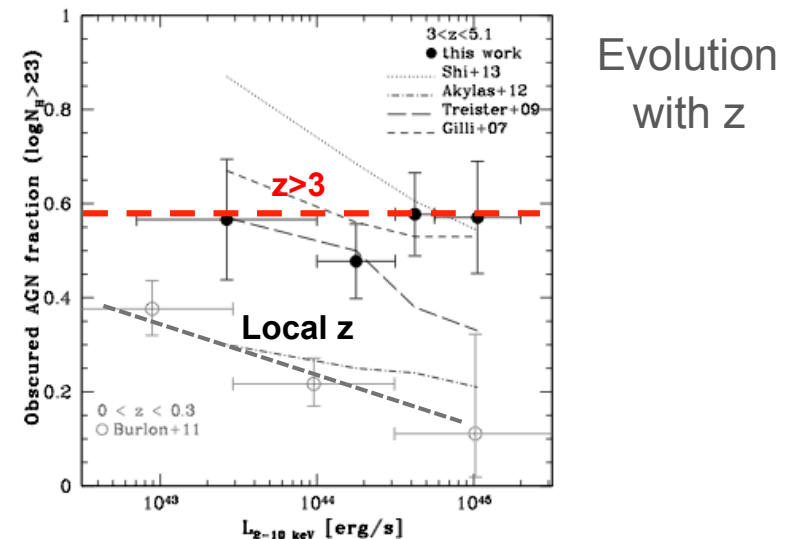
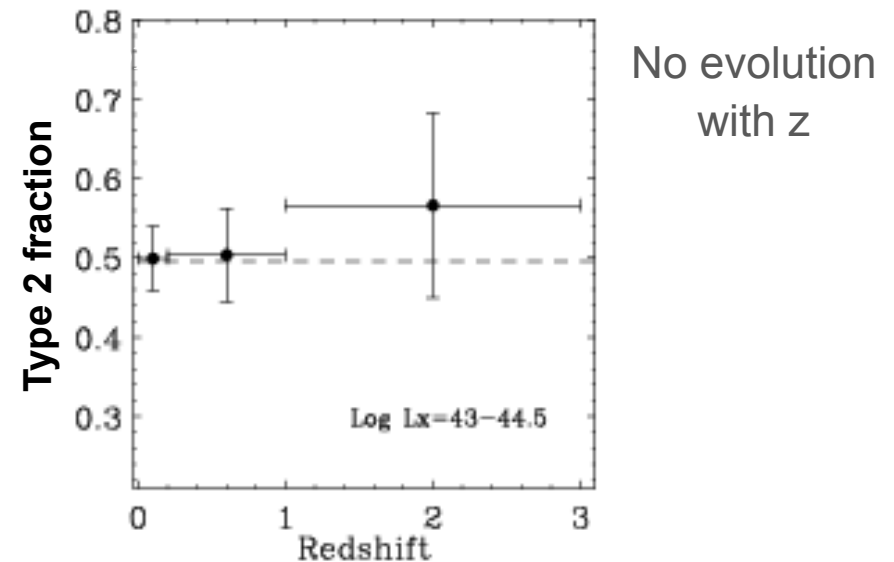


Dependence of the obscured AGN fraction on X-ray luminosity and redshift



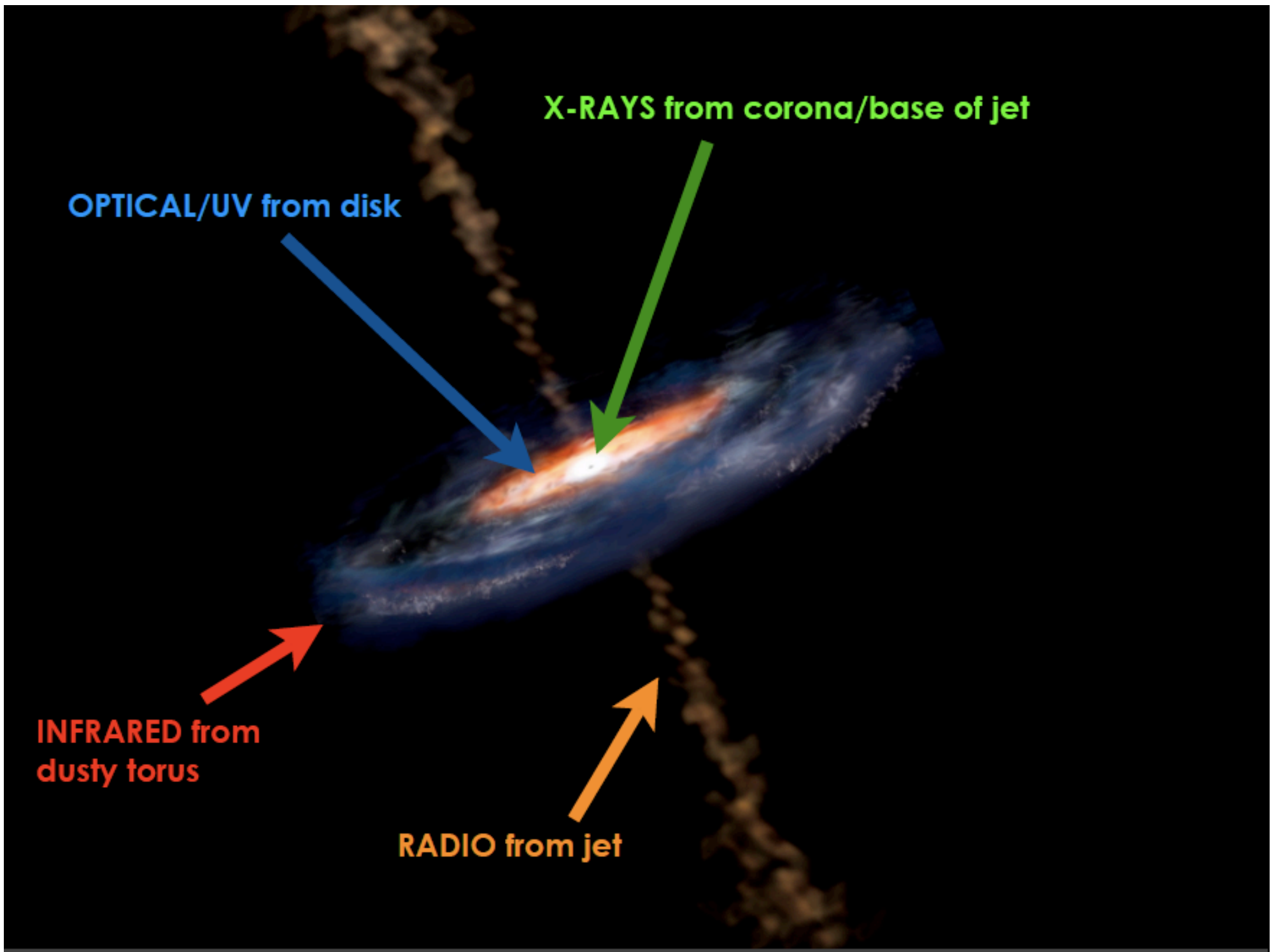
Broad consensus for an obscured AGN fraction declining towards high intrinsic luminosities → *receding torus model* (Lawrence 1991, Simpson 2005; see also Lusso et al. 2013)

Behavior with z still debated (see e.g. La Franca et al. 2005; Treister & Urry 2009; Iwasawa et al. 2012; Vito+13, 14)

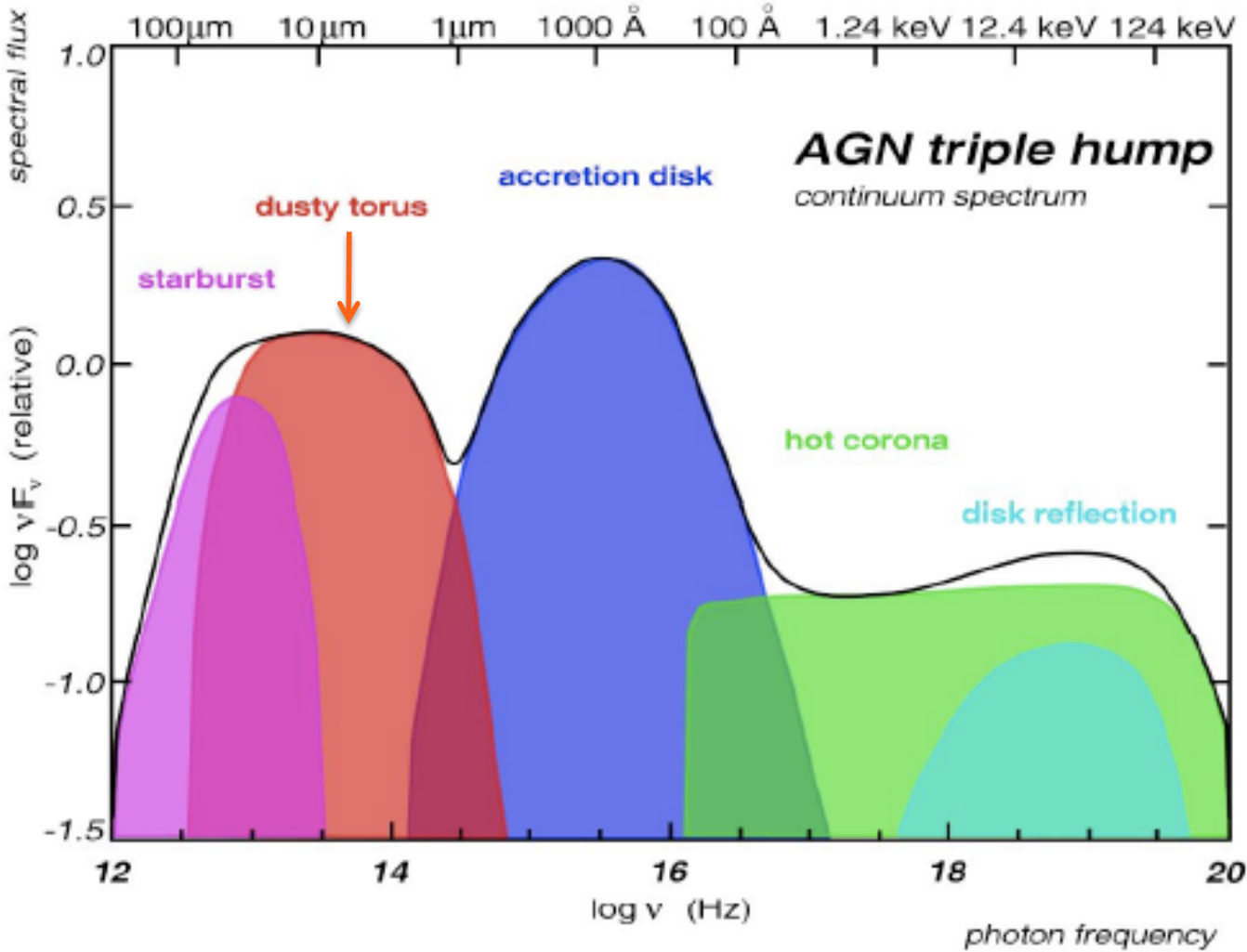


AGN Spectral Energy Distributions.

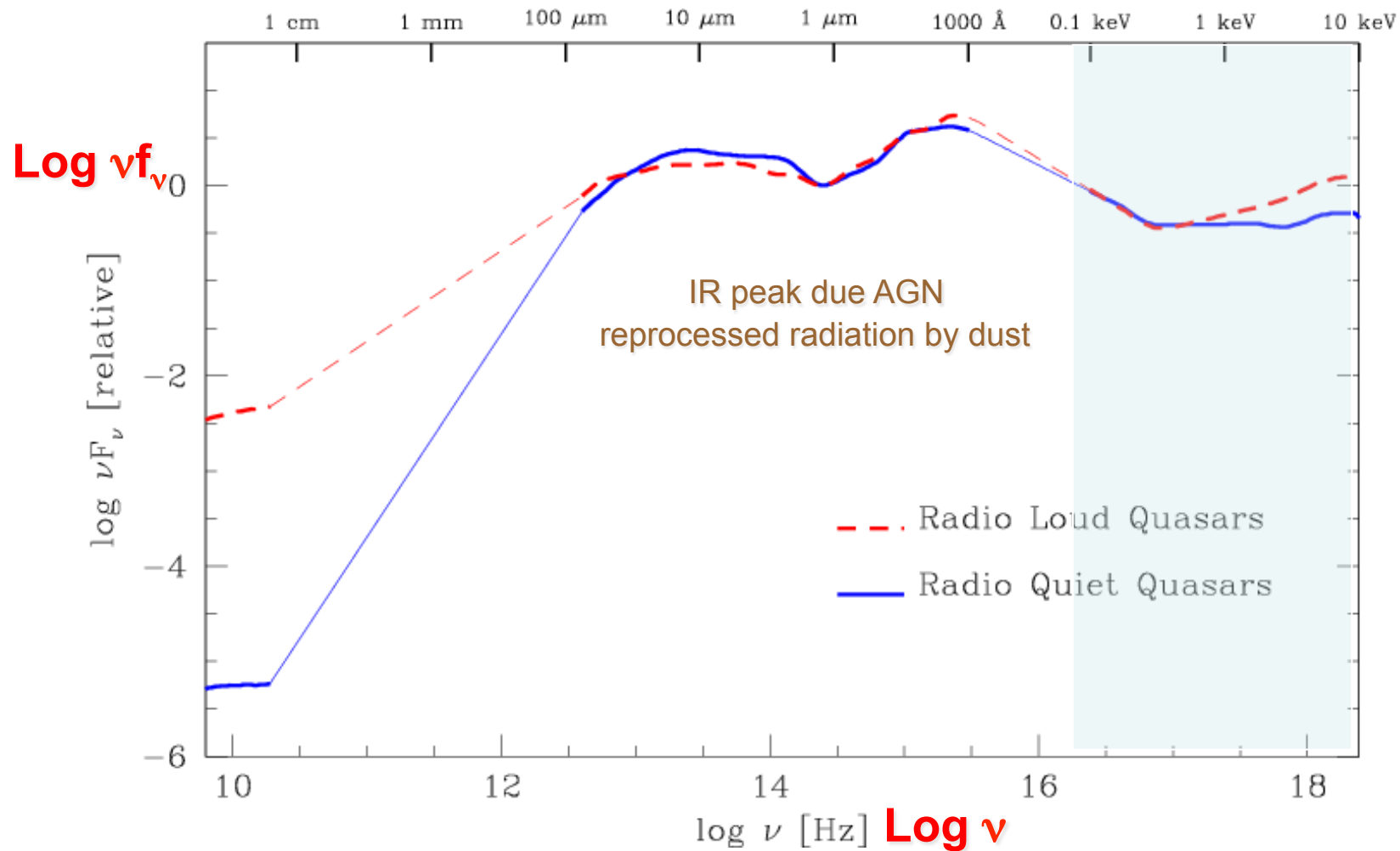
On the properties, location and structure of the X-ray
absorber



Broad-band spectral energy distribution of AGN (I)



Broad-band spectral energy distribution of AGN (II)

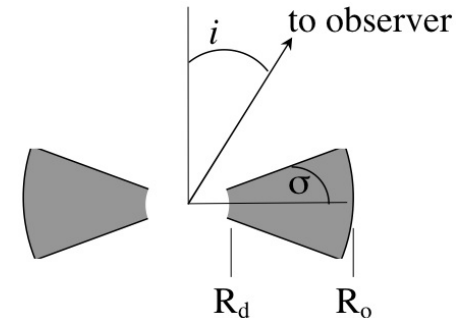


Elvis et al. 1994

Models for the infrared emission of AGN (II)

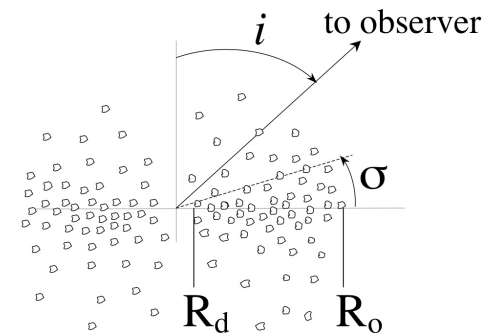
Smooth dust distribution: main properties

- The source is obscured if radiation intercepts the torus, hence obscuration is related to geometrical issues
- Dust temperature is a function of the distance from the source of the radiation field



Clumpy models: main properties

- The probability of direct viewing of the AGN decreases away from the axis, but is always finite
- Different dust temperatures coexist at the same distance from the radiation source, and the same dust temperature occurs at different distances



AGN type is a viewing-dependent probability

Alternative modeling: hydromagnetic disk wind

- Torus=toroidal region of a wind, structured in outflowing clouds. The acceleration is provided by magnetic field lines anchored in the disc (Blandford & Payne '82; Elitzur '08)

Indications from X-ray observations of local Seyferts

Eclipses of the X-ray source are
COMMON in nearby AGN:
 $\Delta N_H \sim 10^{23} - 10^{24} \text{ cm}^{-2}$

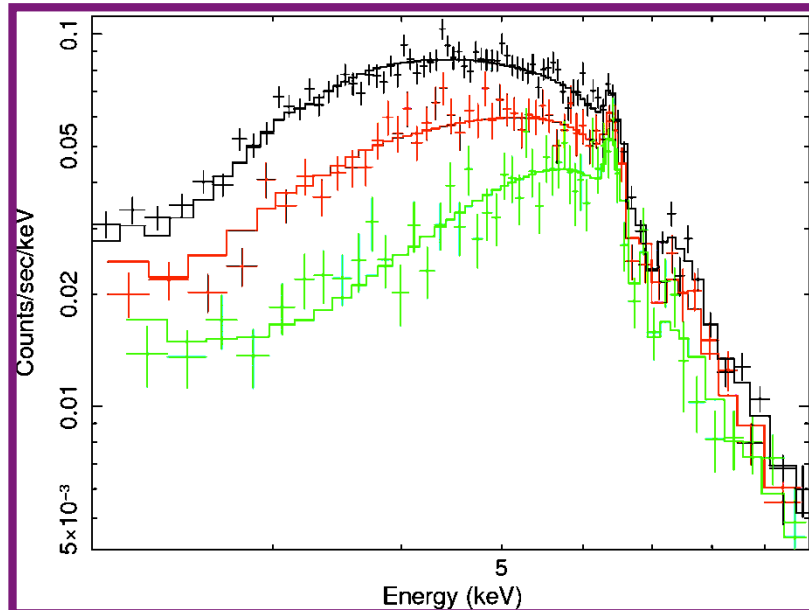
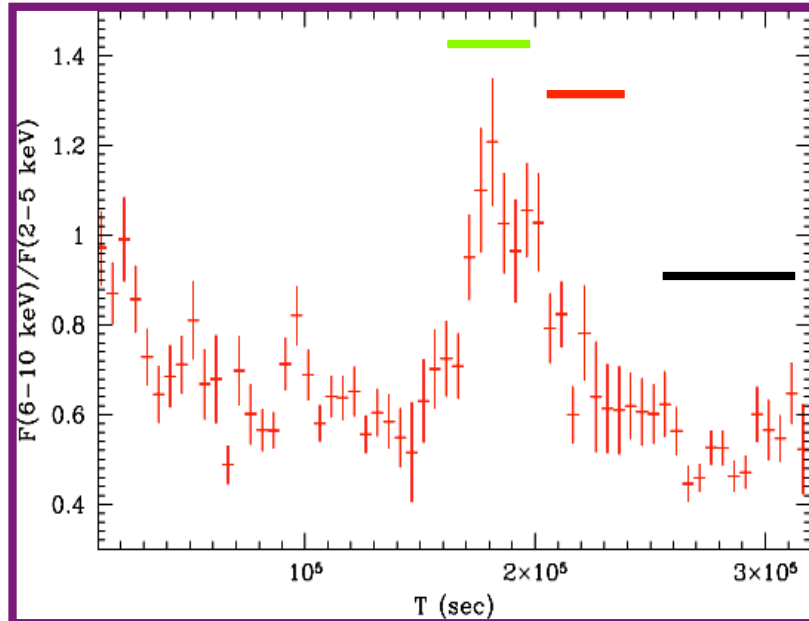


size X-ray src $< 10^{14} \text{ cm}$
 $D < 10^{16} \text{ cm}$

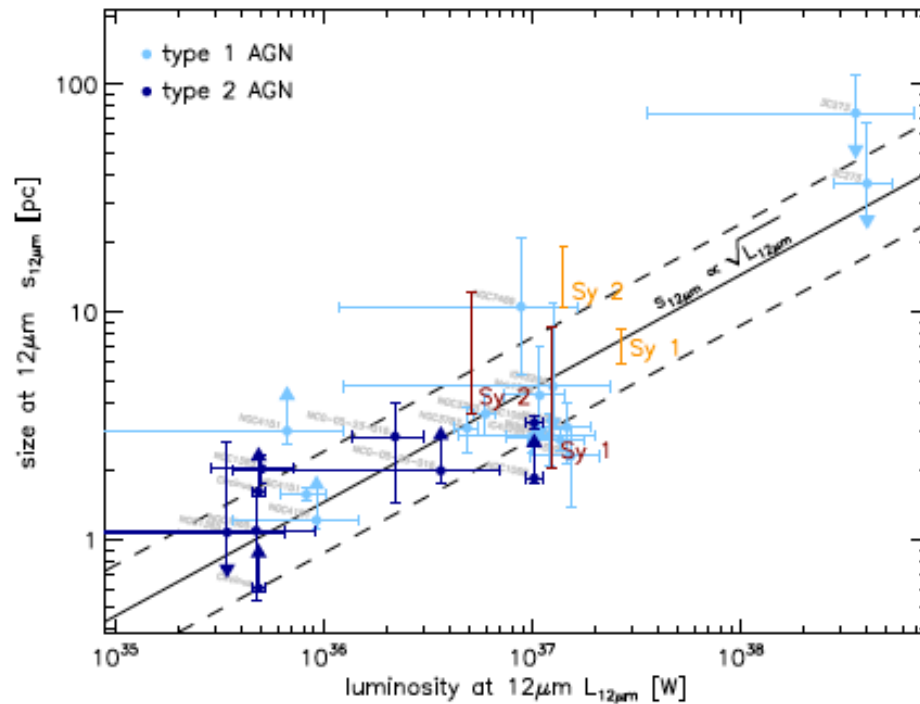


X-ray absorber “made” of BLR
clouds on scales $< \text{pc-scale}$ (torus)

Risaliti et al., 2007, 2010...



High-resolution mid-IR observations of Seyferts

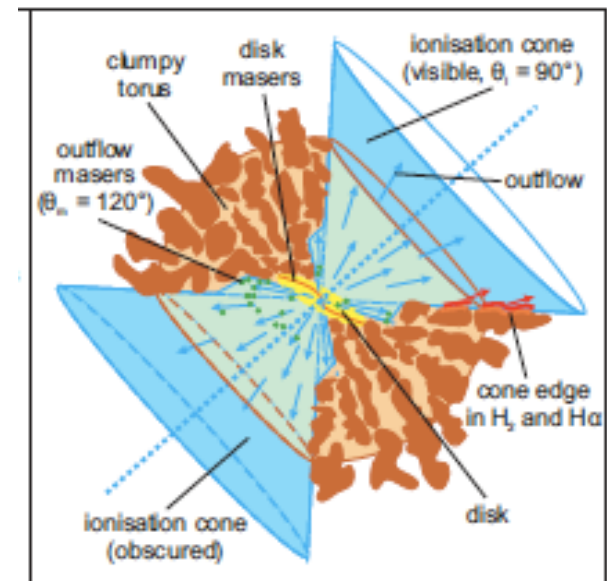


Tristram & Schartmann 2011
 (see also Jaffe+04; Meisenheimer+07;
 Tristram+07; Tristram+09, Burtscher+13)

Tristram+07 - Circinus

- Compact (a few pc) tori with a clumpy/filamentary dust distribution (warm disk + geom. thick torus)

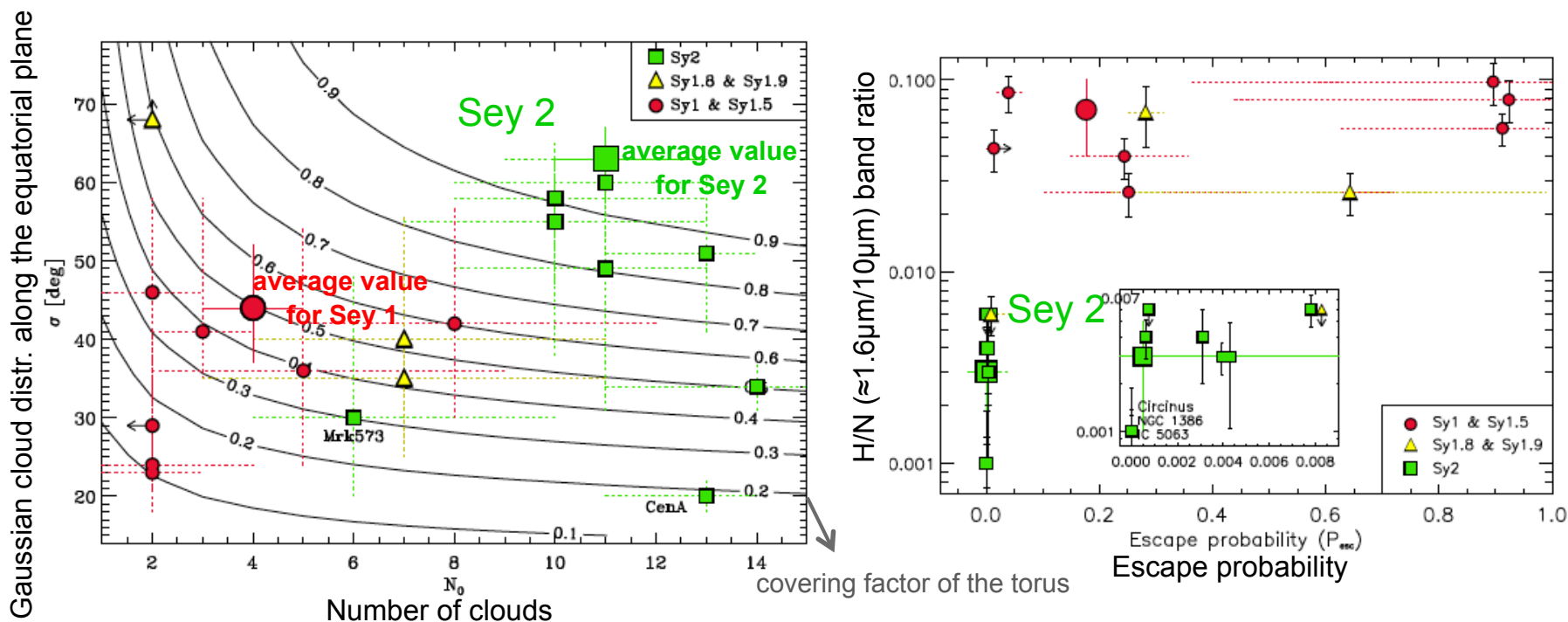
- No significant Sey1/Sey2 difference



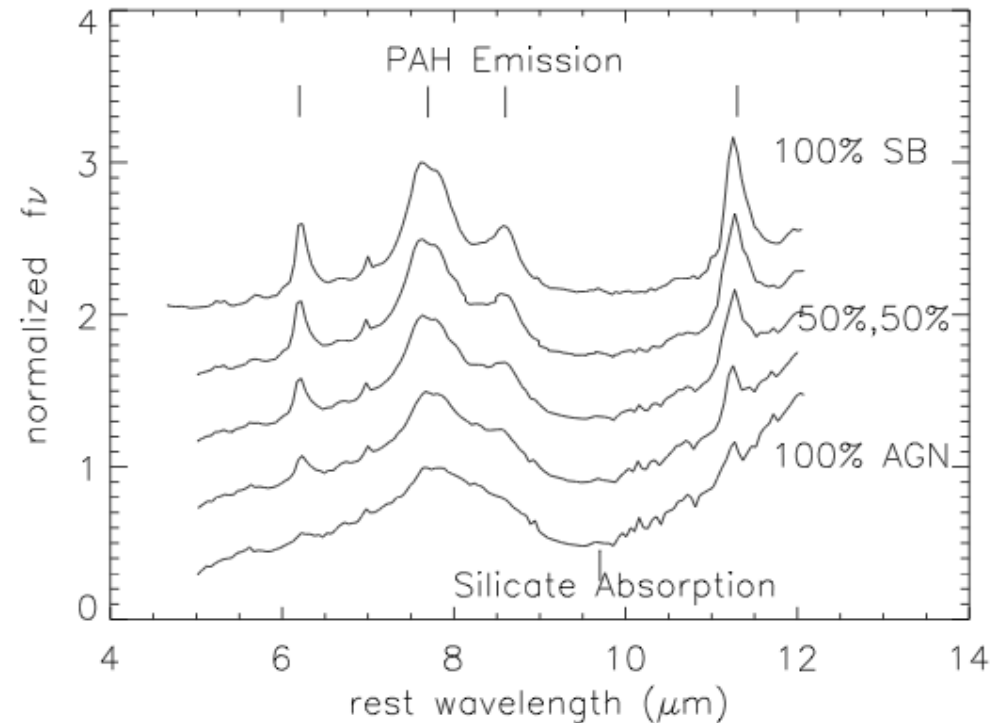
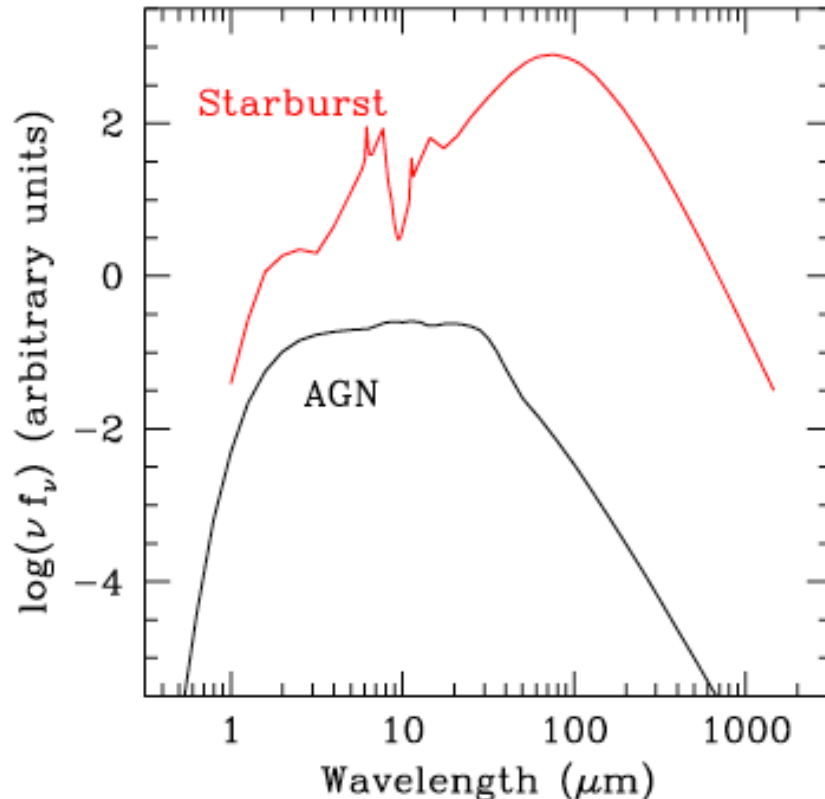
Modeling the mid-IR emission with “clumpy” torus

- ✓ Type 1 vs. Type 2 AGN difference: it is a function of the number of clouds along the line of sight, i.e., of the escape probability
- ✓ Same dust temperatures can be observed at different distances from the AGN

➔ Type 2 AGN: larger number of clouds and lower P_{esc} for the photons to escape



SED fitting: stellar vs. accretion processes

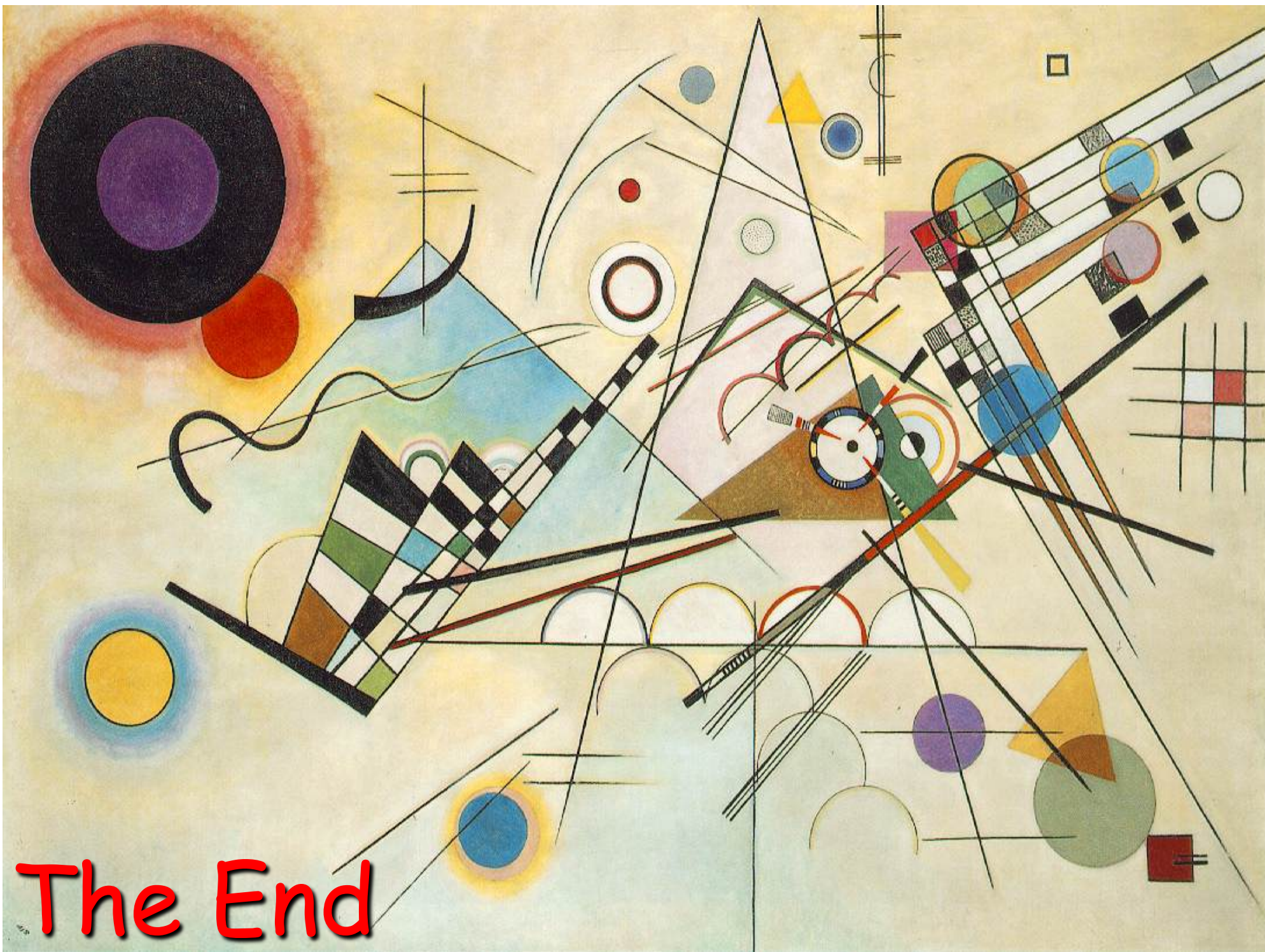


BROAD-BAND SED fitting: common problem to all torus models:
Need to separate the galaxy contribution from that due to the AGN

AGN reprocessed emission and starburst SED peak at different wavelengths

MID-IR continuum vs. PAH features

→ need to decouple the activity due to accretion from that related to stellar processes



The End